

# **Lecture # 21: Stream & Potential Functions for Basic Flows – Part 3**

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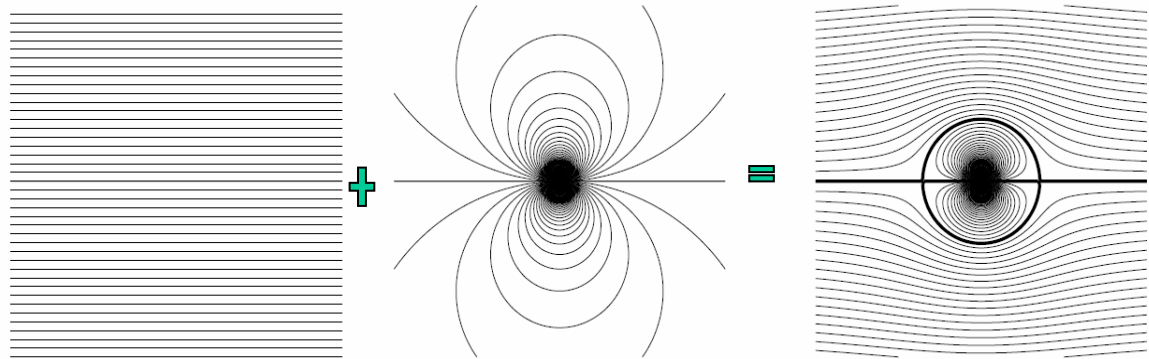
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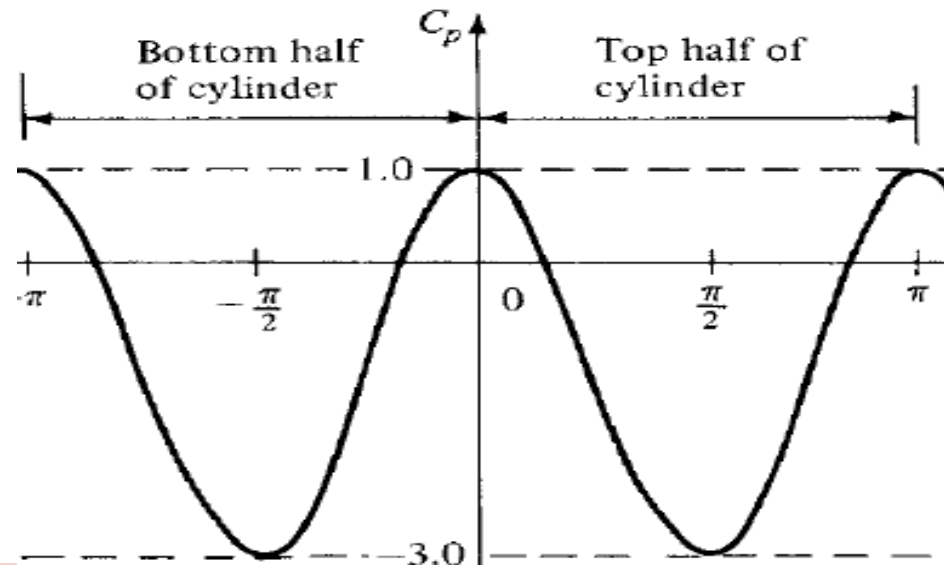
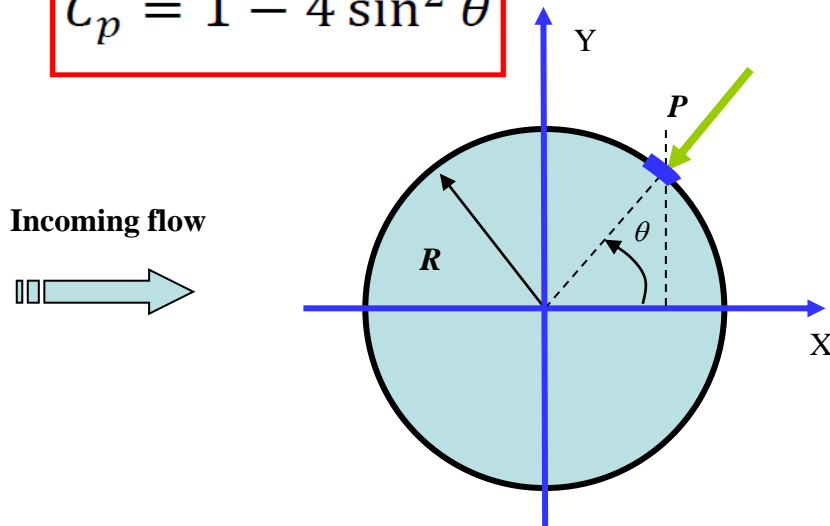
# □ Potential & Stream Functions for Basic Flows

- *Uniform Flow to the Right + A 2-D Doublet*

$$\left. \begin{aligned} \phi &= V_{\infty} r \cos \theta \left( 1 + \frac{R^2}{r^2} \right) \\ \psi &= V_{\infty} r \sin \theta \left( 1 - \frac{R^2}{r^2} \right) \\ V_r &= V_{\infty} \cos \theta \left( 1 - \frac{R^2}{r^2} \right) \\ V_{\theta} &= -V_{\infty} \sin \theta \left( 1 + \frac{R^2}{r^2} \right) \end{aligned} \right\} (r \geq R)$$



$$C_p = 1 - 4 \sin^2 \theta$$

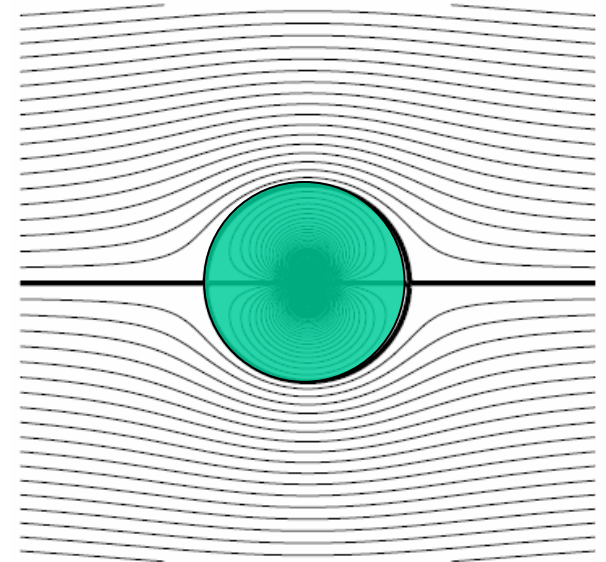


# □ Potential & Stream Functions for Basic Flows

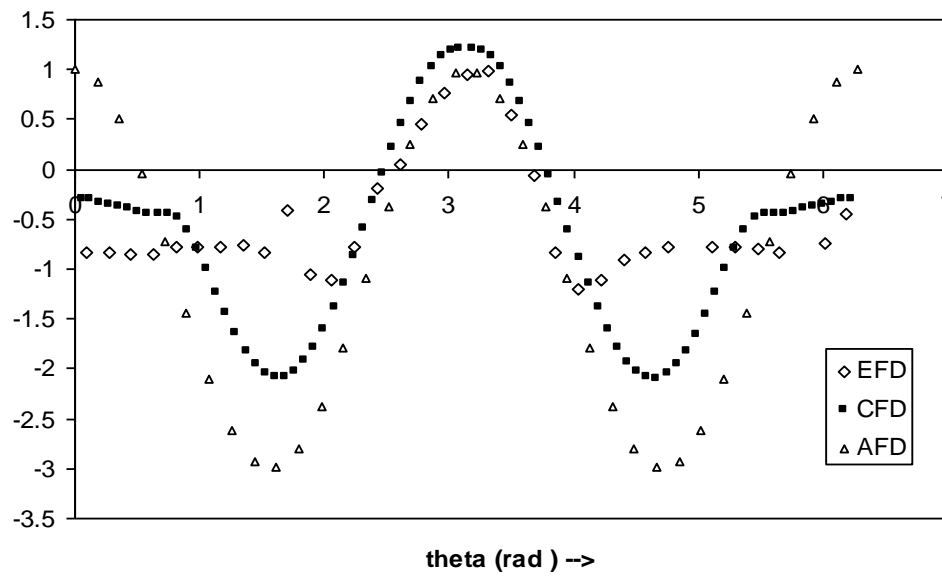
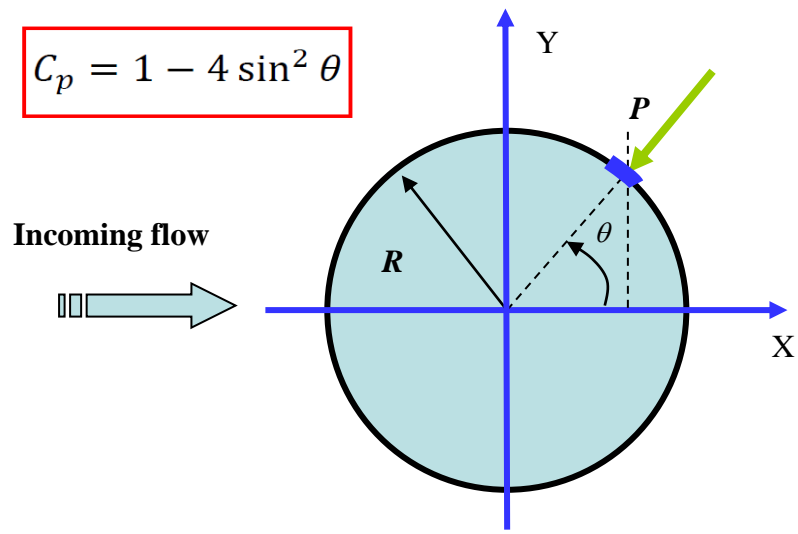
Lift coefficient:  $C_L \equiv \frac{L}{q_\infty S}$

Drag coefficient:  $C_D \equiv \frac{D}{q_\infty S}$

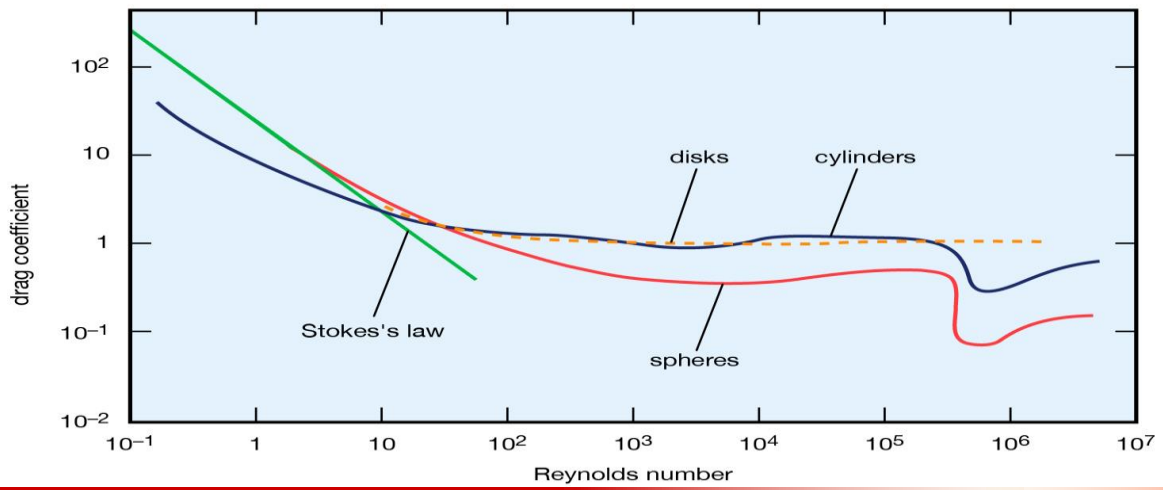
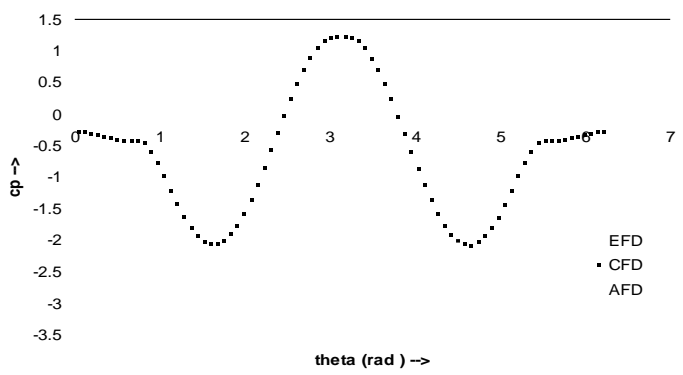
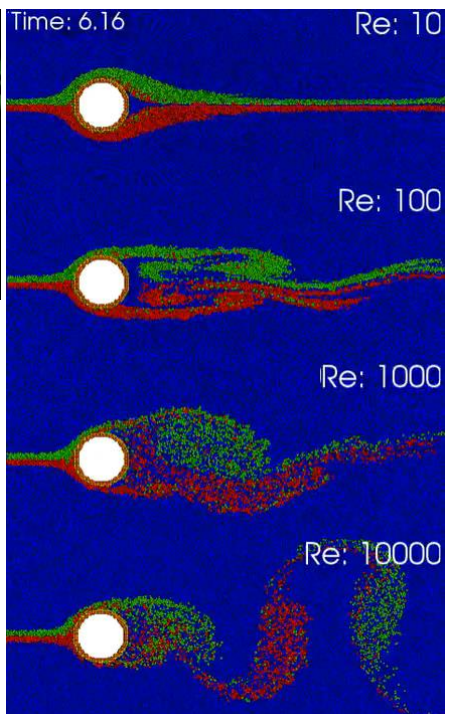
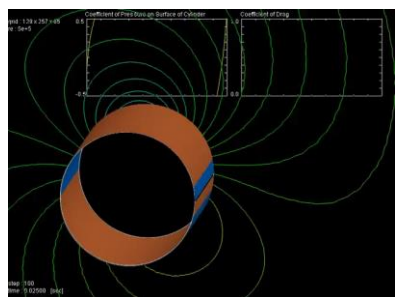
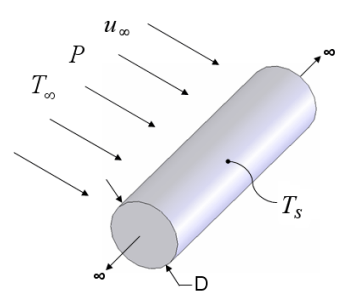
Moment coefficient:  $C_M \equiv \frac{M}{q_\infty S l}$



$C_p = 1 - 4 \sin^2 \theta$



# DRAG COEFFICIENT OF A CIRCULAR CYLINDER IN A REAL FLOW



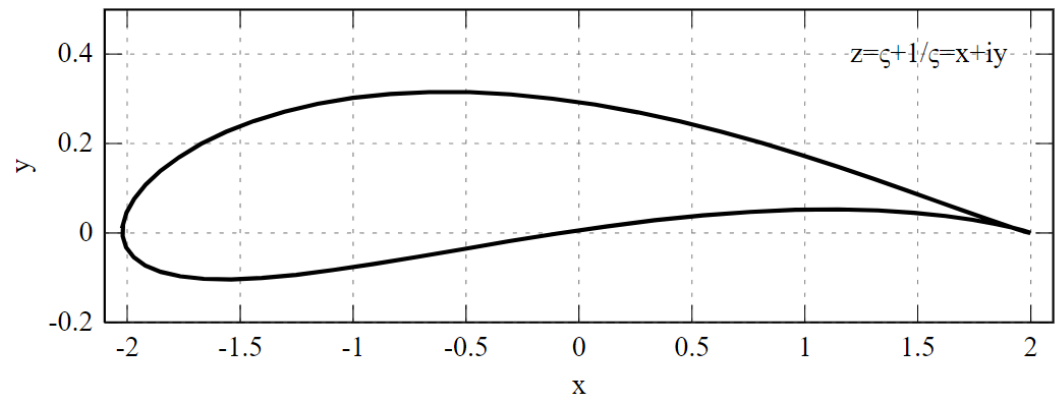
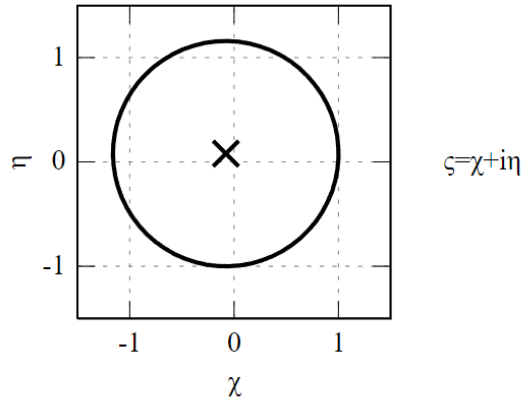
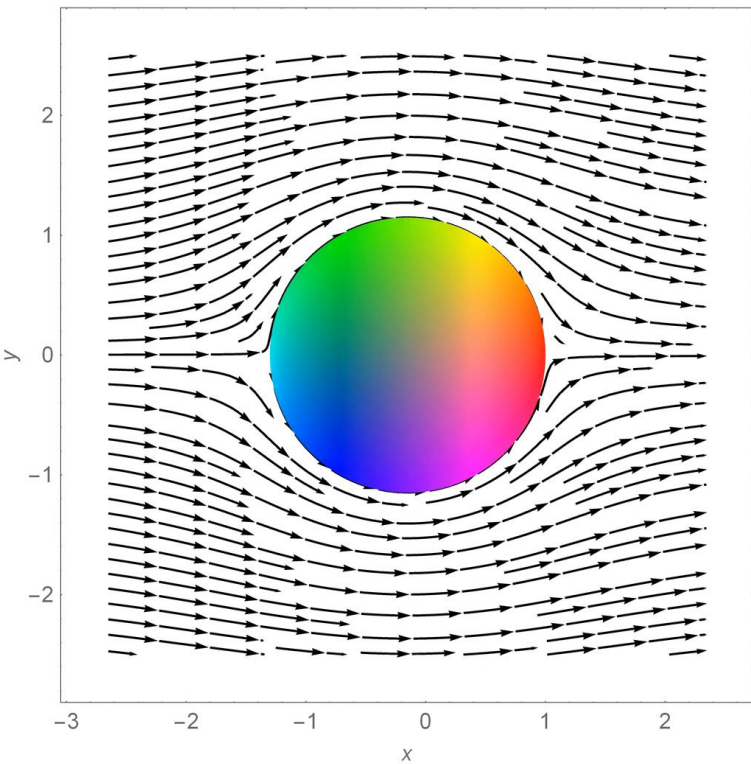
Shape	Drag Coefficient
Sphere	0.47
Half-sphere	0.42
Cone	0.50
Cube	1.05
Angled Cube	0.80
Long Cylinder	0.82
Short Cylinder	1.15
Streamlined Body	0.04
Streamlined Half-body	0.09

Measured Drag Coefficients

# □ Zhukovsky TRANSFORM – TOPICS TO BE COVERED IN AERE541

- *By using Zhukovsky transformation, the spinning circle above can be transformed into the Zoukovsky airfoil below.*

Zhukovsky transformation of potential flow over a cylinder



# □ Potential & Stream Functions for Basic Flows

## • 2D Vortex flow

A 2-D point vortex is a mathematical concept that induces a velocity field given by

$$V_r = 0, \quad V_\theta = \frac{\text{const.}}{r} = \frac{C}{r}$$

1. Check if the flow satisfies conservation of mass (Is it a physically possible flow?)

$$\nabla \cdot \vec{V} \stackrel{?}{=} 0$$

$$\nabla \cdot \vec{V} = \frac{1}{r} \left[ \frac{\partial(V_r r)}{\partial r} + \frac{\partial V_\theta}{\partial \theta} \right] = 0 \rightarrow \psi \text{ exist.}$$

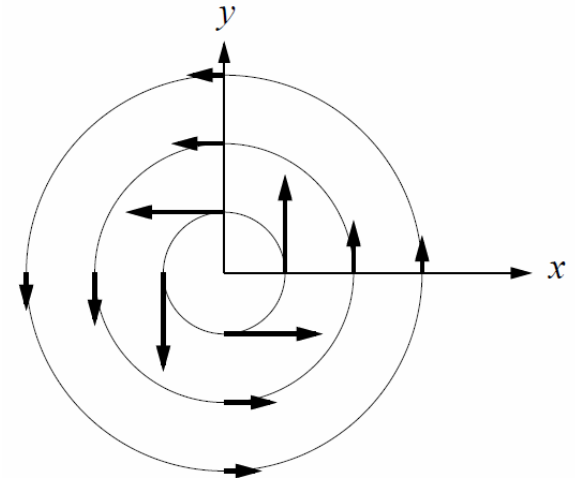
$$V_r = \frac{\partial \psi}{r \partial \theta} = 0 \rightarrow \psi = g(r)$$

$$V_\theta = -\frac{\partial \psi}{\partial r} = \frac{C}{r} \rightarrow \psi = -C \ln r + f(\theta)$$

$$\frac{\partial \psi}{\partial \theta} = f'(\theta) = 0$$

$$f(\theta) = \text{const.}$$

$$\therefore \psi = -C \ln r + \text{const.}$$



When  $r \rightarrow 0$ ,  $V_\theta = \infty$  and  $\psi \rightarrow \infty$ . To eliminate the infinite velocity it is arbitrary assumed that  $\psi = 0$  at  $r = R$

$$\therefore \psi = -C \ln R + \text{const.} = 0$$

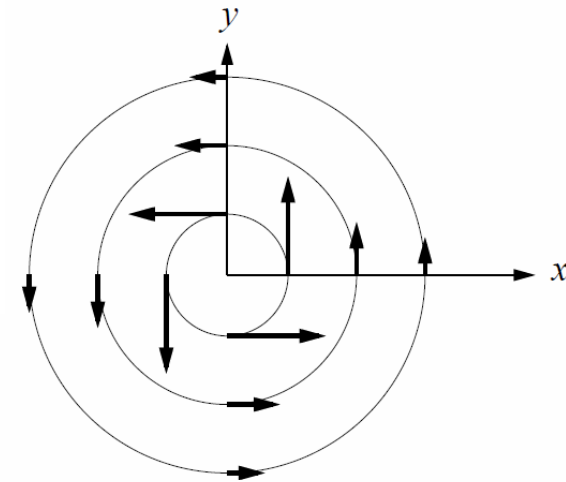
$$\text{const.} = C \ln R$$

$$\psi = -C \ln \left( \frac{r}{R} \right) \quad \text{for } (r \geq R)$$

# □ Potential & Stream Functions for Basic Flows

## • 2D Vortex Flow

$$V_r = 0, \quad V_\theta = \frac{\text{const.}}{r} = \frac{C}{r}$$



Check if the flow is irrotational

$$\nabla \times \vec{V} \stackrel{?}{=} 0$$

$$\nabla \times \vec{V} = \frac{1}{h_1 h_2 h_3} \begin{vmatrix} h_1 \hat{e}_1 & h_2 \hat{e}_2 & h_3 \hat{e}_3 \\ \frac{\partial}{\partial q_1} & \frac{\partial}{\partial q_2} & \frac{\partial}{\partial q_3} \\ h_1 V_1 & h_2 V_2 & h_3 V_3 \end{vmatrix} = \frac{1}{r} \begin{vmatrix} \hat{e}_r & r \hat{e}_\theta & \hat{e}_z \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial z} \\ V_r & r V_\theta & V_z \end{vmatrix}$$

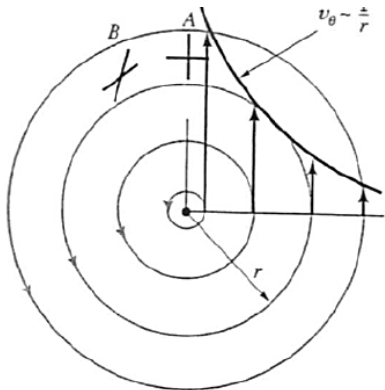
Vorticity or  $(\nabla \times \vec{V})$  in the  $r - \theta$  plane

$$\frac{1}{r} \left( \frac{\partial r V_\theta}{\partial r} - \frac{\partial V_r}{\partial \theta} \right) = \left( \frac{\partial C}{\partial r} - \frac{\partial 0}{\partial \theta} \right) = 0 \rightarrow \phi \text{ exist.}$$

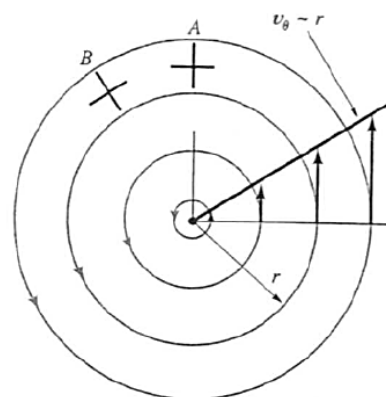
# □ Potential & Stream Functions for Basic Flows

## Free and forced vortices

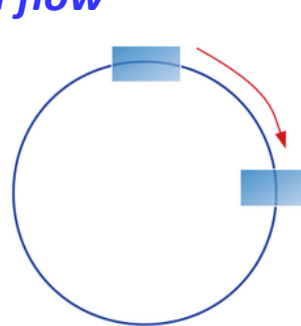
- Free or potential vortex is irrotational :  $v_{\theta}r = const.$
- Forced vortex is rotational:  $\frac{v_{\theta}}{r} = const.$



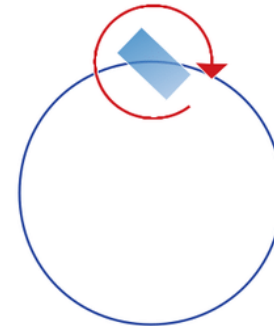
• **Irrotational flow**



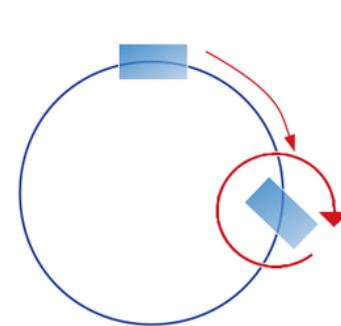
• **Rotational flow**



Circulation



Rotation



Circulation &  
Rotation



# □ Potential & Stream Functions for Basic Flows

- The velocity components are

$$v_r = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = 0$$

$$v_\theta = -\frac{\partial \psi}{\partial r} = \frac{c_1}{r}$$

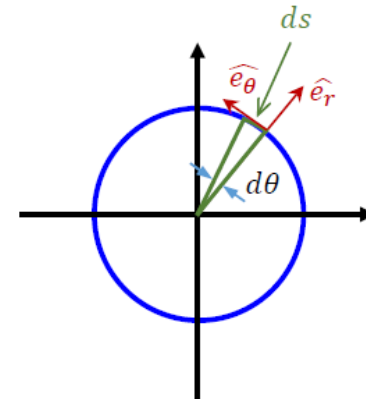
- Now let's define circulation  $\Gamma$  around a circle with radius  $r$ .

Adding velocity contribution from each term

$$\Gamma = \oint \vec{V} \cdot d\vec{s}$$

Note:

$$\begin{aligned}\vec{V} &= v_\theta \hat{e}_\theta \\ d\vec{s} &= r d\theta \hat{e}_\theta\end{aligned}$$



Therefore:

$$\Gamma = \oint v_\theta \hat{e}_\theta \cdot r d\theta \hat{e}_\theta = \int_0^{2\pi} \frac{c_1}{r} r d\theta = c_1 \int_0^{2\pi} d\theta = 2\pi c_1 \rightarrow c_1 = \frac{\Gamma}{2\pi}$$

# □ Potential & Stream Functions for Basic Flows

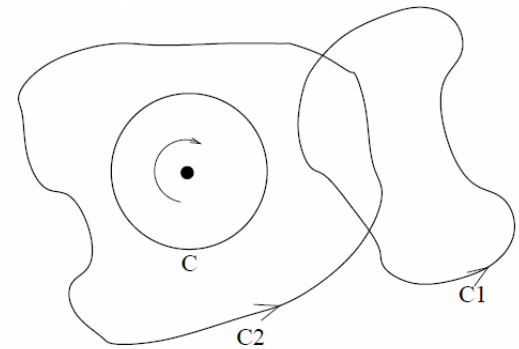
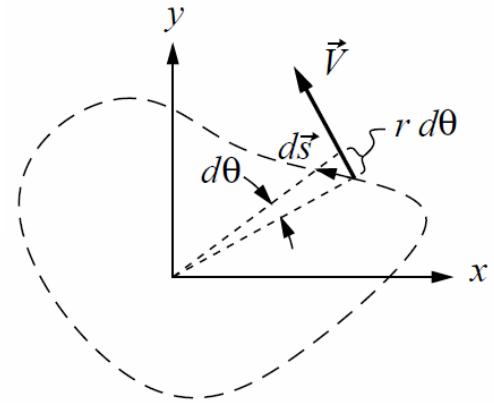
Around closed curve  $C1$  that does not include the point vortex

$$\Gamma_{C1} = - \oint_{C1} \vec{V} \cdot d\vec{l} = \iint_{S1} \underbrace{(\nabla \times \vec{V})}_0 \cdot d\vec{A} = 0$$

Around  $C2$  that includes the point vortex.

$$\begin{aligned} \Gamma_{C2} &= - \left[ \oint_{C2} (V_r \hat{e}_r + V_\theta \hat{e}_\theta) \cdot (dr \hat{e}_r + r d\theta \hat{e}_\theta) \right] \\ &= - \left[ \oint_C (V_r \hat{e}_r + V_\theta \hat{e}_\theta) \cdot (dr \hat{e}_r + r d\theta \hat{e}_\theta) \right] + \left[ \oint_{C2-C} (V_r \hat{e}_r + V_\theta \hat{e}_\theta) \cdot (dr \hat{e}_r + r d\theta \hat{e}_\theta) \right] \\ &= - \left[ \oint_C (V_r \hat{e}_r + V_\theta \hat{e}_\theta) \cdot (dr \hat{e}_r + r d\theta \hat{e}_\theta) + 0 \right] \\ &= - \left[ \oint_C V_r dr + \oint_C V_\theta r d\theta \right] = - \left[ 0 + \int_0^{2\pi} \left( \frac{C}{r} \right) r d\theta \right] = -2\pi C \end{aligned}$$

$$\Gamma_{C2} = -2\pi C \quad \text{or} \quad -\frac{\Gamma}{2\pi} = C$$



This implies that the circulation evaluated for a curve enclosing the 2-D vortex is a constant and not equal to zero

$$\Gamma = \begin{cases} -2\pi C & , \text{ (circuit encircles origin)} \\ 0 & , \text{ (circuit doesn't encircle origin)} \end{cases}$$

# □ Potential & Stream Functions for Basic Flows

- The final form is:

$$\psi = -\frac{\Gamma}{2\pi} \ln r$$
$$v_r = 0, \quad v_\theta = \frac{\Gamma}{2\pi r}$$

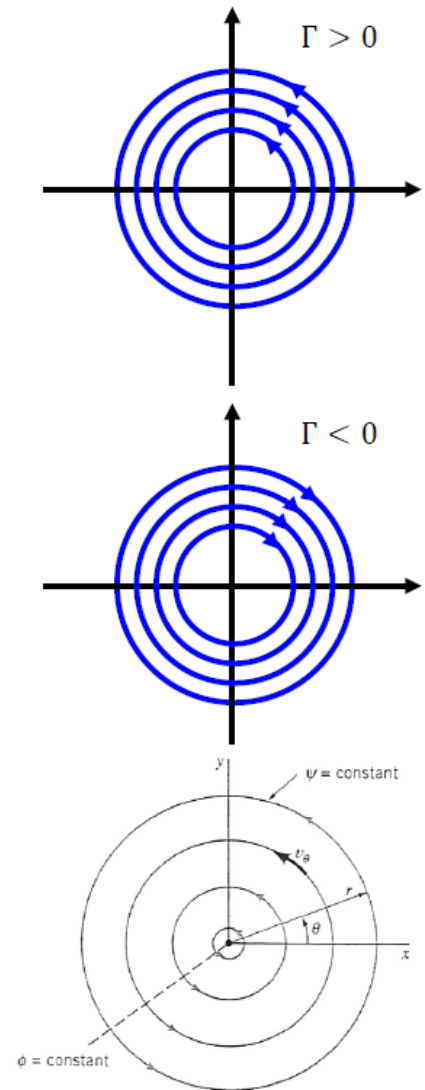
Potential function

$$v_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = \frac{\Gamma}{2\pi r} \rightarrow \phi(r, \theta) = \int \frac{\Gamma}{2\pi} d\theta$$

$$\rightarrow \phi = \frac{\Gamma}{2\pi} \theta + f(r)$$

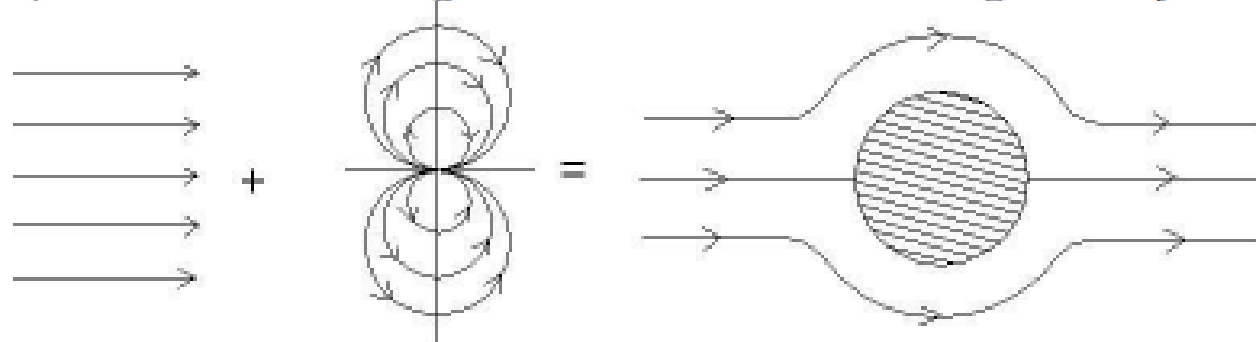
$$v_r = \frac{\partial \phi}{\partial r} = 0 \rightarrow f(r) = c = 0$$

$$\phi = \frac{\Gamma}{2\pi} \theta$$

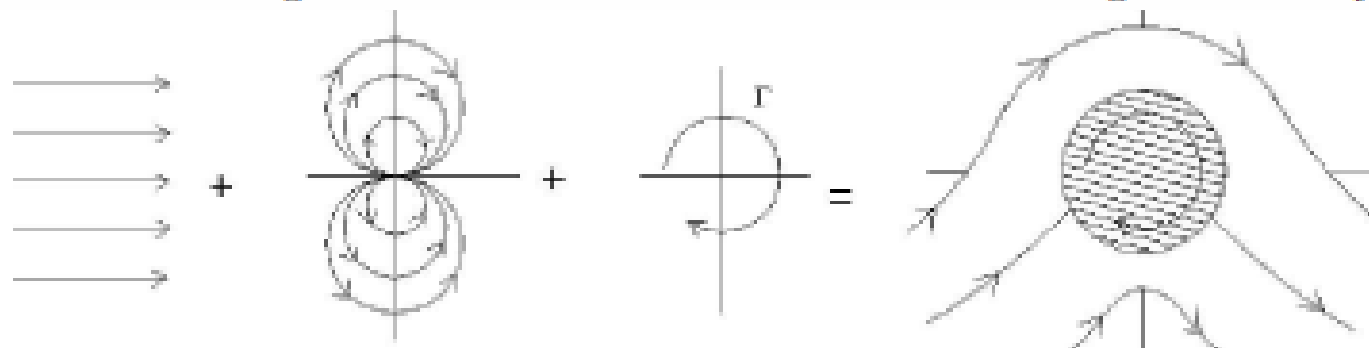


# □ Potential & Stream Functions for Basic Flows

- As we all know, uniform flow to the right + 2-D Doublet = non-lifting over a cylinder



- Uniform flow to the right + 2-D Doublet + 2-D Point Vortex = Lifting flow over a cylinder



The parameters for lifting flow over a cylinder are as follow (spinning cylinder):

Quantity	Non-lifting flow over a cylinder	Vortex of Strength $\Gamma$	Combination
$\psi$	$V_{\infty} r \sin \theta (1 - \frac{R^2}{r^2})$	$\frac{\Gamma}{2\pi} \ln \frac{r}{R}$	$V_{\infty} r \sin \theta (1 - \frac{R^2}{r^2}) + \frac{\Gamma}{2\pi} \ln \frac{r}{R}$
$\phi$	$V_{\infty} r \cos \theta (1 + \frac{R^2}{r^2})$	$-\frac{\Gamma}{2\pi} \theta$	$V_{\infty} r \cos \theta (1 + \frac{R^2}{r^2}) - \frac{\Gamma}{2\pi} \theta$
$V_r$	$V_{\infty} \cos \theta (1 - \frac{R^2}{r^2})$	0	$V_{\infty} \cos \theta (1 - \frac{R^2}{r^2})$
$V_{\theta}$	$-V_{\infty} \sin \theta (1 + \frac{R^2}{r^2})$	$-\frac{\Gamma}{2\pi r}$	$-V_{\infty} \sin \theta (1 + \frac{R^2}{r^2}) - \frac{\Gamma}{2\pi r}$

# □ Potential & Stream Functions for Basic Flows

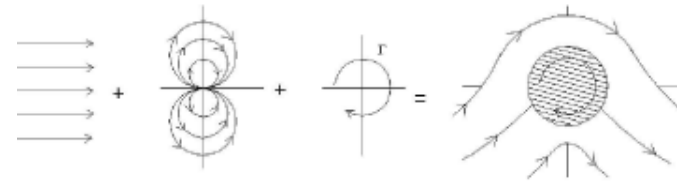
Superposition of free vortex and non-lifting cylinder flow

- Let's add a vortex flow to the non-lifting flow around the cylinder, i.e.

$$\psi = \psi_{\text{uniform}} + \psi_{\text{doublet}} + \psi_{\text{vortex}}$$

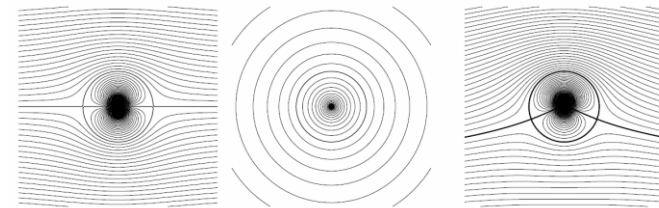
Recall for cylinder flow:

$$\psi = V_{\infty} r \sin \theta \left( 1 - \frac{R^2}{r^2} \right)$$



Therefore:

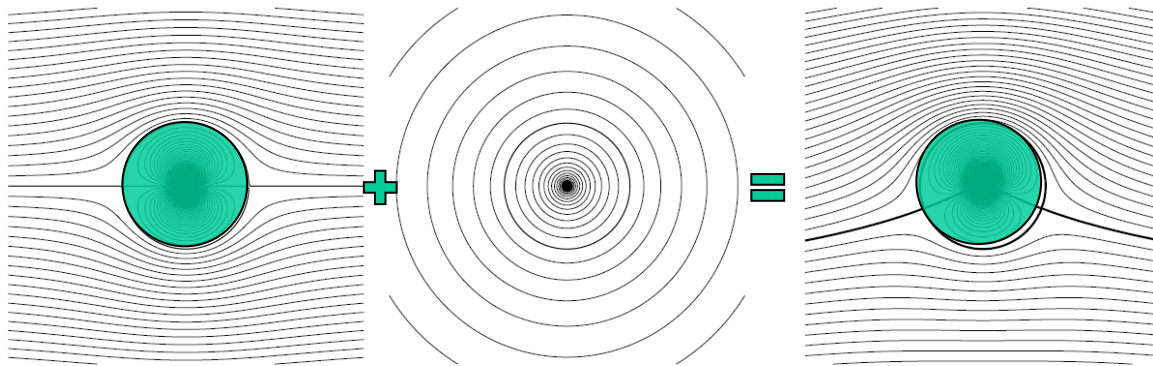
$$\psi = V_{\infty} r \sin \theta \left( 1 - \frac{R^2}{r^2} \right) - \frac{\Gamma}{2\pi} \ln r + c$$



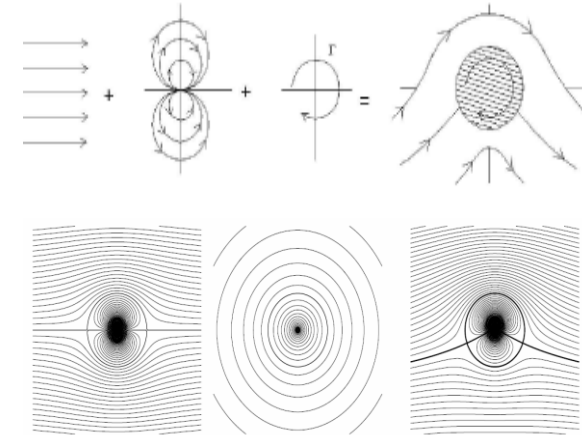
We can always add a constant to stream function. Adding  $c$  here to keep  $\psi = 0$  on the surface:

$$r = R \rightarrow \psi = 0 = 0 - \frac{\Gamma}{2\pi} \ln R + c \rightarrow c = \frac{\Gamma}{2\pi} \ln R$$

# □ Potential & Stream Functions for Basic Flows



Lifting flow past a cylinder



$$\psi = V_{\infty} r \sin \theta \left( 1 - \frac{R^2}{r^2} \right) + \frac{\Gamma}{2\pi} \ln \frac{R}{r}$$

Velocity components

$$v_r = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = V_{\infty} \cos \theta \left( 1 - \frac{R^2}{r^2} \right)$$

$$v_{\theta} = -\frac{\partial \psi}{\partial r} = -V_{\infty} \sin \theta - V_{\infty} \sin \theta \frac{R^2}{r^2} + \frac{\Gamma}{2\pi r}$$

At  $r = R$

$$v_r = 0, v_{\theta} = -2V_{\infty} \sin \theta + \frac{\Gamma}{2\pi R}$$

# □ Potential & Stream Functions for Basic Flows

Stagnation points are located at  $v_r = 0, v_\theta = 0$

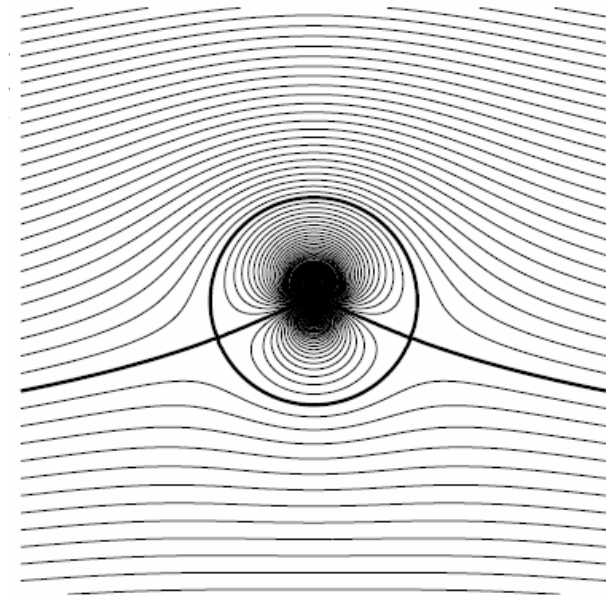
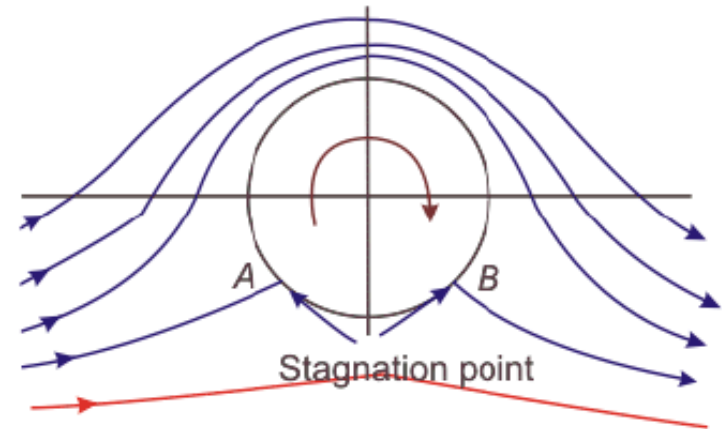
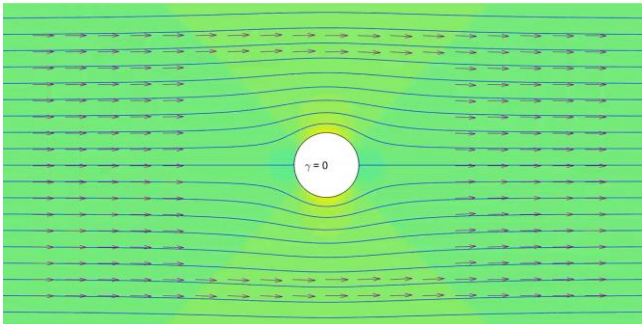
$$v_\theta = -2V_\infty \sin \theta + \frac{\Gamma}{2\pi R} = 0$$

$$\sin \theta = \frac{\Gamma}{4\pi V_\infty R}$$

On cylinder surface  $y = R \sin \theta = \frac{\Gamma}{4\pi V_\infty}$

Note  $y > 0$  for  $\Gamma > 0$  and  $y < 0$  for  $\Gamma < 0$

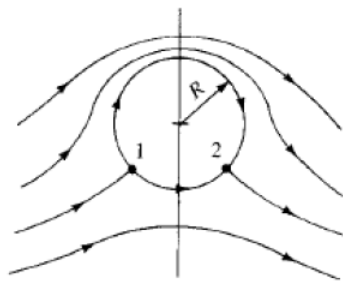
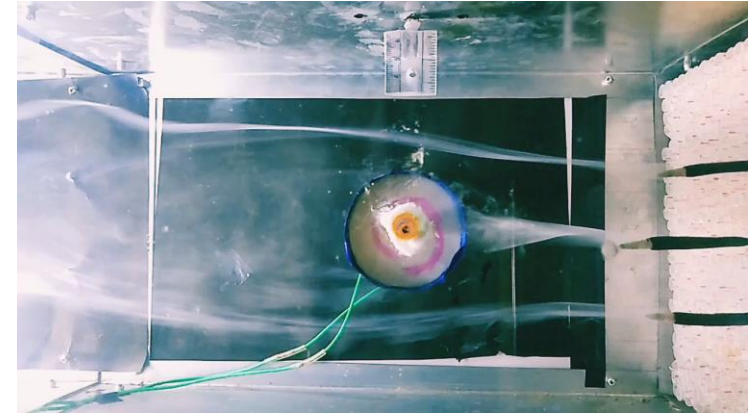
(remember we defined  $\Gamma > 0$  counter-clockwise but textbook is backward, so let's consider  $\Gamma < 0$ !)



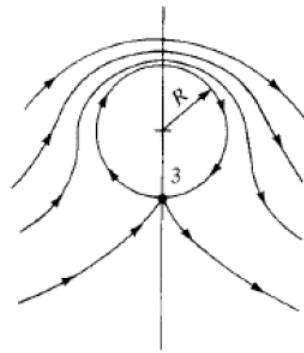
# □ Potential & Stream Functions for Basic Flows

$$v_r = \frac{1}{r} \frac{\partial \psi}{\partial \theta} = V_\infty \cos \theta \left( 1 - \frac{R^2}{r^2} \right)$$

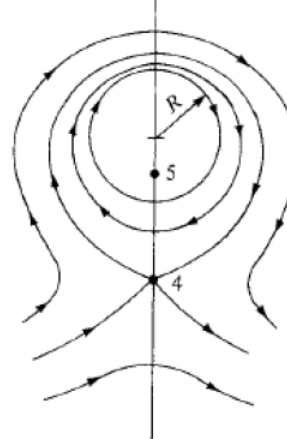
$$v_\theta = -\frac{\partial \psi}{\partial r} = -V_\infty \sin \theta - V_\infty \sin \theta \frac{R^2}{r^2} + \frac{\Gamma}{2\pi r}$$



$$|\Gamma| < 4\pi V_\infty R$$



$$|\Gamma| = 4\pi V_\infty R$$



$$|\Gamma| > 4\pi V_\infty R$$

For  $|\Gamma| > 4\pi V_\infty R$ , the equation for  $v_\theta = 0$  doesn't give us a solution. We need to look at the  $v_r = 0$

$$v_r = V_\infty \cos \theta \left( 1 - \frac{R^2}{r^2} \right) = 0 \rightarrow r = R \text{ or } \cos \theta = 0 \rightarrow \theta = \pm \frac{\pi}{2}$$

$$r = \frac{\Gamma}{4\pi V_\infty} \pm \sqrt{\left( \frac{\Gamma}{4\pi V_\infty} \right)^2 - R^2}$$



# □ Potential & Stream Functions for Basic Flows

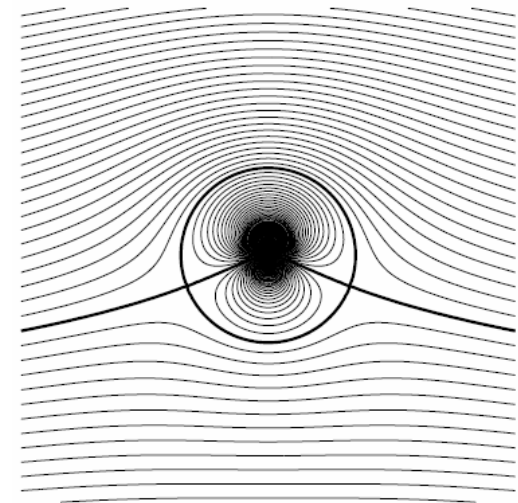
Pressure coefficient

$$V^2 = v_r^2 + v_\theta^2 = 0 + \left(-2V_\infty \sin \theta + \frac{\Gamma}{2\pi R}\right)^2$$
$$C_p = 1 - \left(\frac{V}{V_\infty}\right)^2$$

$$C_p = 1 - \left(-2 \sin \theta + \frac{\Gamma}{2\pi R V_\infty}\right)^2$$

$$C_p = 1 - \left[4 \sin^2 \theta - \frac{2\Gamma \sin \theta}{\pi R V_\infty} + \left(\frac{\Gamma}{2\pi R V_\infty}\right)^2\right]$$

$$p - p_\infty = \frac{1}{2} \rho V_\infty^2 C_p$$



# □ Potential & Stream Functions for Basic Flows

## Pressure integral

- Force on a small arc on the cylinder

$$dF = (p - p_\infty)lds = (p - p_\infty)lRd\theta$$

$$F = \int_0^{2\pi} dF$$

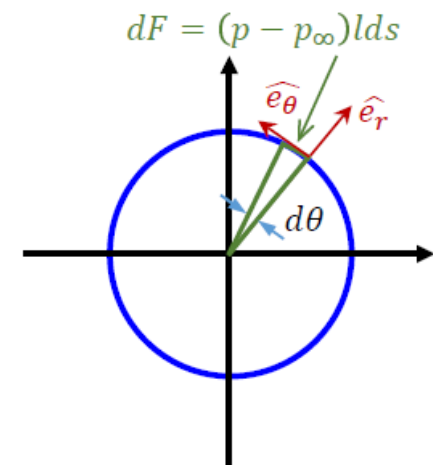
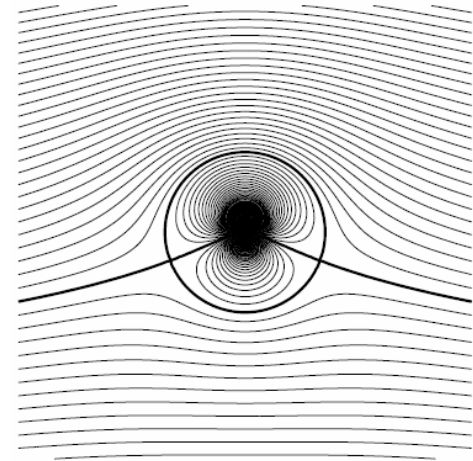
Or in terms of its components:

$$D = F_x = F \cos \theta = \int_0^{2\pi} (p - p_\infty)lR \cos \theta d\theta$$

$$L = F_y = F \sin \theta = \int_0^{2\pi} (p - p_\infty)lR \sin \theta d\theta$$

$$C_D = \frac{D}{\frac{1}{2}\rho V_\infty^2 2Rl} = \frac{\int_0^{2\pi} (p - p_\infty)lR \cos \theta d\theta}{\frac{1}{2}\rho V_\infty^2 2Rl}$$

$$= \frac{1}{2} \int_0^{2\pi} \frac{(p - p_\infty)}{\frac{1}{2}\rho V_\infty^2} \cos \theta d\theta = \frac{1}{2} \int_0^{2\pi} C_p \cos \theta d\theta$$



Note we are using the projection area  $2Rl$  as the reference area for  $C_D$

# □ Potential & Stream Functions for Basic Flows

Drag on the cylinder

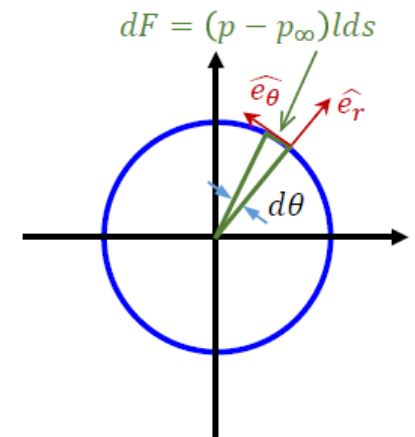
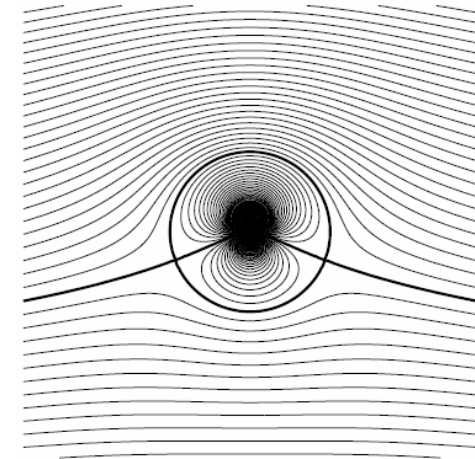
$$\begin{aligned} C_D &= \frac{1}{2} \int_0^{2\pi} C_p \cos \theta \, d\theta \\ &= \frac{1}{2} \int_0^{2\pi} \left( \cos \theta \right. \\ &\quad \left. - \left[ 4 \sin^2 \theta \cos \theta - \frac{2\Gamma \sin \theta \cos \theta}{\pi R V_\infty} + \left( \frac{\Gamma}{2\pi R V_\infty} \right)^2 \cos \theta \right] \right) d\theta \end{aligned}$$

Note  $\int_0^{2\pi} \cos \theta \, d\theta = 0$ ,  $\int_0^{2\pi} \sin^2 \theta \cos \theta \, d\theta = 0$ ,  
 $\int_0^{2\pi} \sin \theta \cos \theta \, d\theta = 0$

Therefore

$$C_D = 0$$

d'Alembert  
paradox



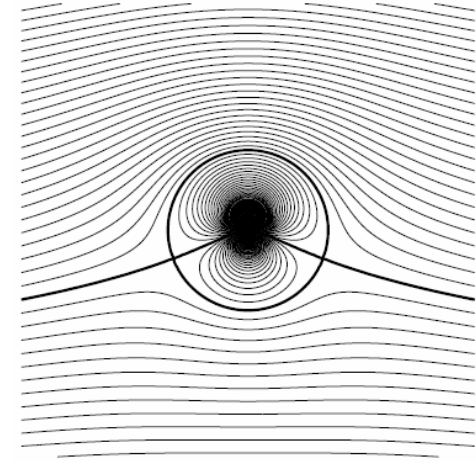
Note we are using the  
projection area  $2Rl$  as  
the reference area for  
 $C_D$

# □ Potential & Stream Functions for Basic Flows

## Lift on the cylinder

$$C_L = \frac{1}{2} \int_0^{2\pi} C_p \sin \theta \, d\theta$$

$$= \frac{1}{2} \int_0^{2\pi} \left( \sin \theta - \left[ 4 \sin^3 \theta - \frac{2\Gamma \sin^2 \theta}{\pi R V_\infty} + \left( \frac{\Gamma}{2\pi R V_\infty} \right)^2 \sin \theta \right] \right) d\theta$$



Note  $\int_0^{2\pi} \sin \theta \, d\theta = 0$ ,  $\int_0^{2\pi} \sin^3 \theta \, d\theta = 0$ ,  $\int_0^{2\pi} \sin^2 \theta \, d\theta = \pi$

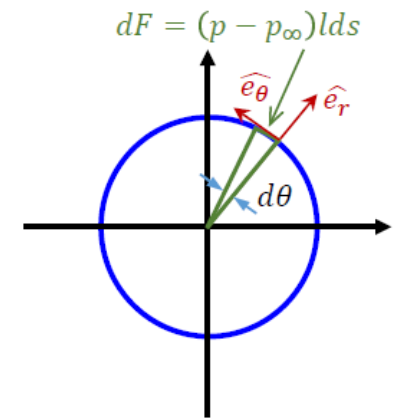
Therefore

$$C_L = \frac{\Gamma}{\pi R V_\infty} (\pi) \rightarrow C_L = \frac{\Gamma}{R V_\infty}$$

$$L = \frac{1}{2} \rho V_\infty^2 (2Rl) \frac{\Gamma}{R V_\infty} = \rho V_\infty \Gamma l$$

$$L' = \rho V_\infty \Gamma$$

Kutta-Joukowski theorem



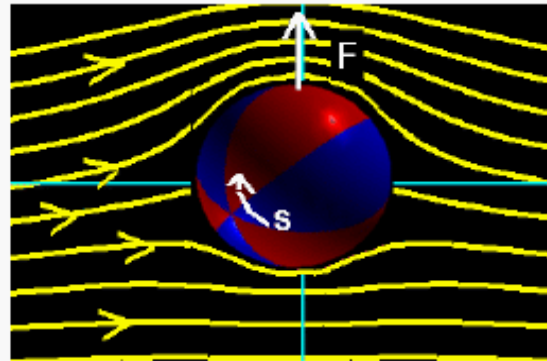
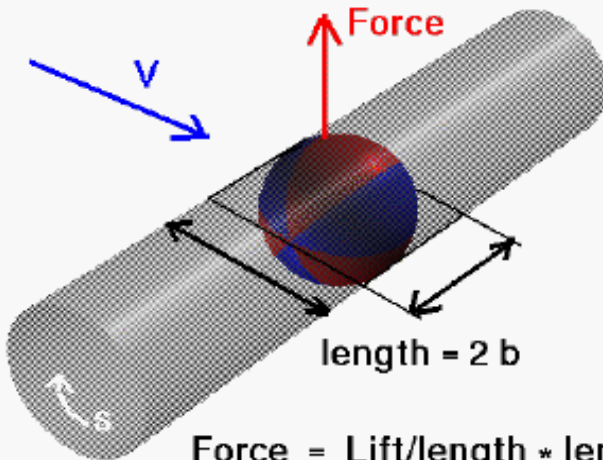
Note we are using the projection area  $2Rl$  as the reference area for  $C_D$

# □ THE KUTTA-JOUKOWSKI LIFT THEOREM

- **The Kutta-Joukowski Lift Theorem** states the lift per unit length of a spinning cylinder is equal to the density ( $\rho$ ) of the air times the strength of the rotation ( $\Gamma$ ) times the velocity ( $V$ ) of the air.

**Kutta-Joukowski Lift Theorem for a Cylinder:**

Lift per unit length of a cylinder acts perpendicular to the velocity ( $V$ ) and is given by:  $L = r G V$  (lbs/ft)

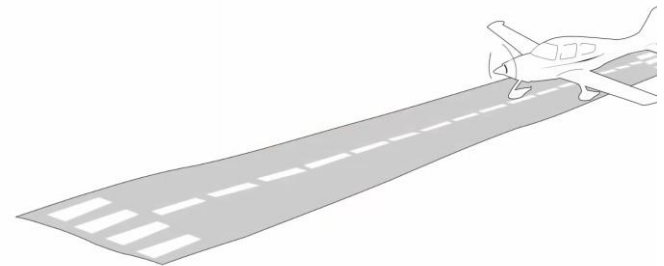


$$L' = \rho V_{\infty} \Gamma$$

Force = Lift/length \* length \* area of circle / area of square

- **Theories in the Production of Lift**

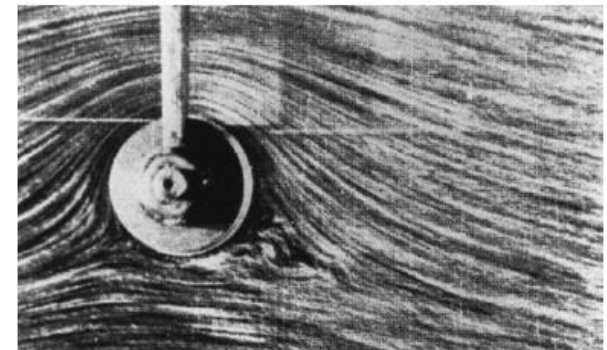
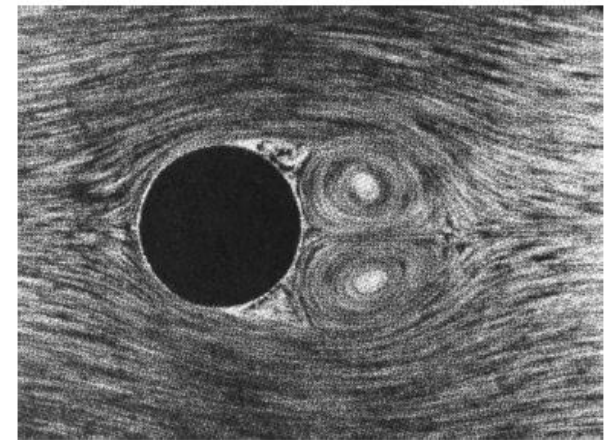
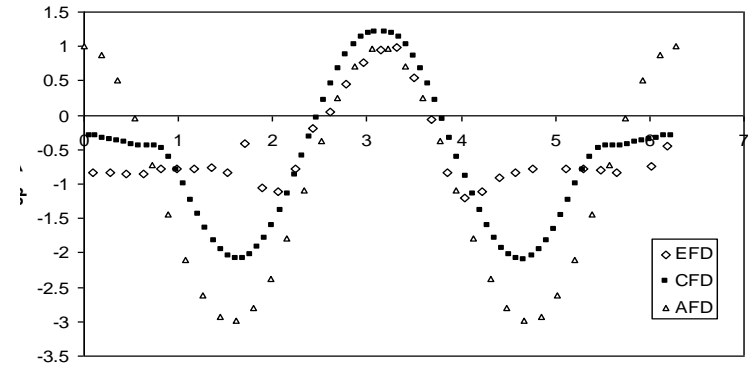
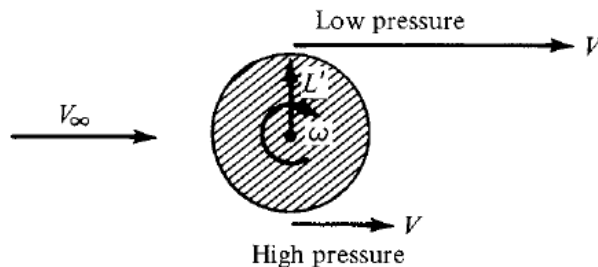
[https://www.youtube.com/watch?v=mBRCUU\\_fQrQ](https://www.youtube.com/watch?v=mBRCUU_fQrQ)



# □ Potential & Stream Functions for Basic Flows

## Theoretical vs real flow

- In real flow past cylinder  $C_D \neq 0$ 
  - Drag is primarily caused by the flow separation, a viscous effect. The potential theory neglects the viscous effects hence can not predict the drag force.
- Flow around a spinning cylinder resembles that of theory and generates lift proportional to the circulation



# □ Potential & Stream Functions for Basic Flows

## Theoretical vs real flows

- For streamlined shapes (airfoils), potential flow can predict the pressure distribution and the resulting forces with a good accuracy
  - Low angle of attacks, where the flow does not separate
  - Can estimate pressure drag (as long as flow is attached)!
  - Can't get friction drag!

