

# **Lecture # 25: Airfoil Aerodynamics - Part 03: Symmetrical Airfoil – 01**

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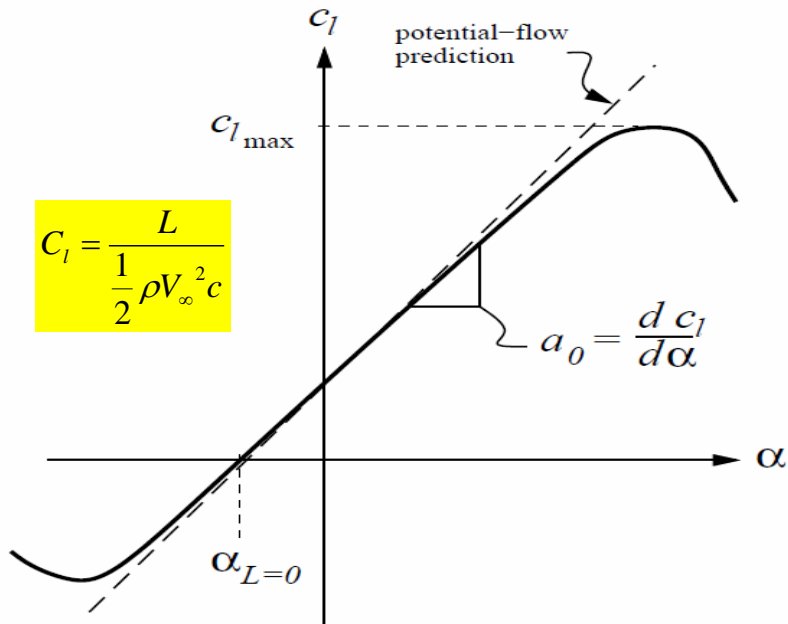
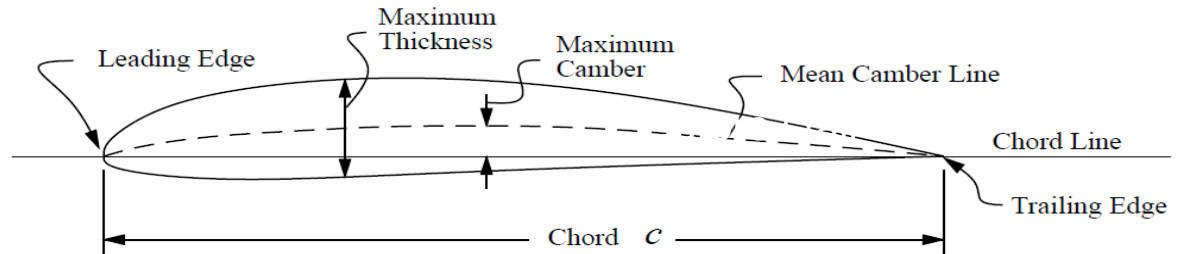
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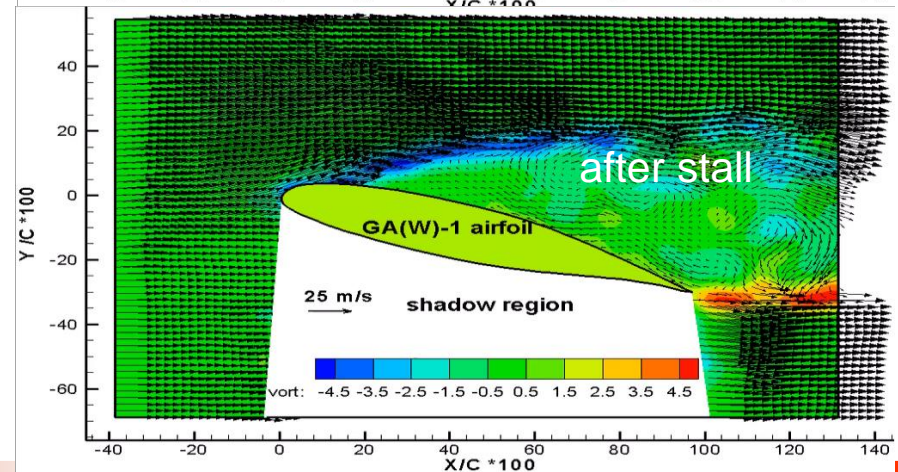
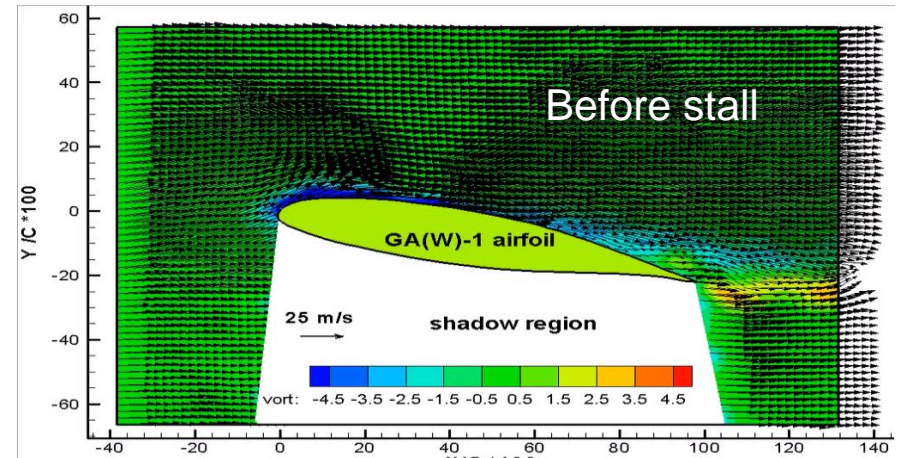
# Thin Airfoil Theory

## Assumptions:

- 2-dimensions
- Inviscid\*
- Incompressible\*
- Irrotational\*
- Small  $\alpha$
- Small  $\max \tau / c$  (i.e., airfoil thickness)
- Small  $\max z / c$  (i.e., airfoil camber)



$$C_l = \frac{L}{\frac{1}{2} \rho V_\infty^2 c}$$

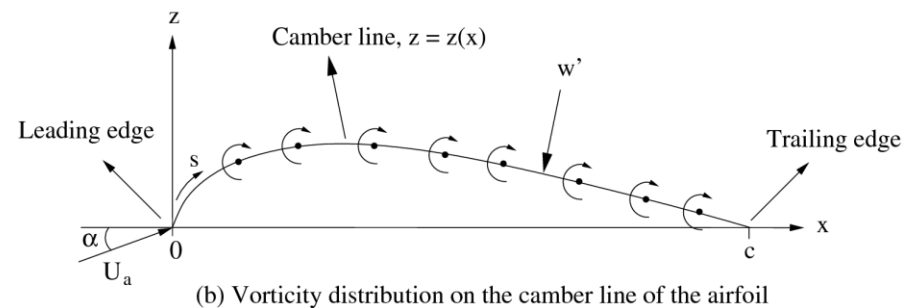
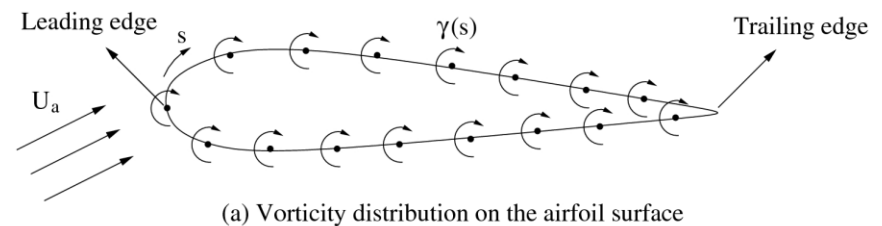
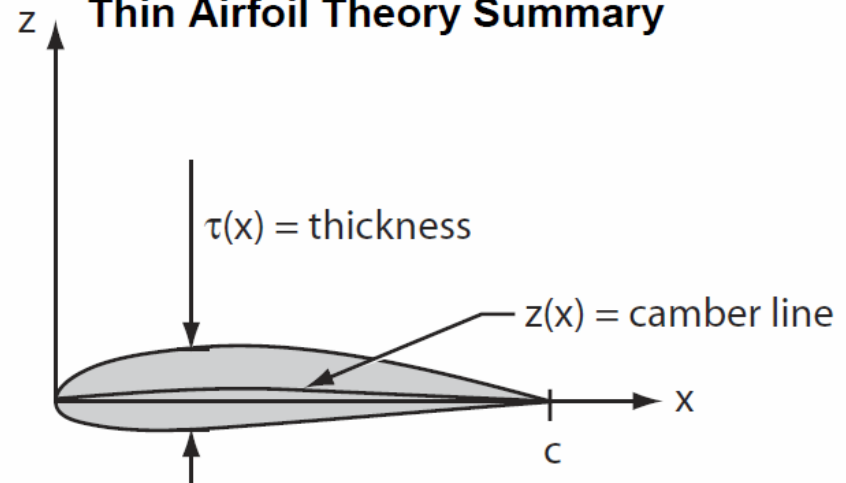


# Thin Airfoil Theory

## Thin Airfoil Approximation

- For a sufficiently thin airfoil we can approximate the shape of the airfoil by its mean camber line.
- In the thin airfoil approximation then, we distribute vortex sheet on its mean camber line instead of top/bottom surfaces.
  - The strength of the vortex is adjusted such that the camber line is a streamline
- This approximation allows a closed form solution to be obtained.

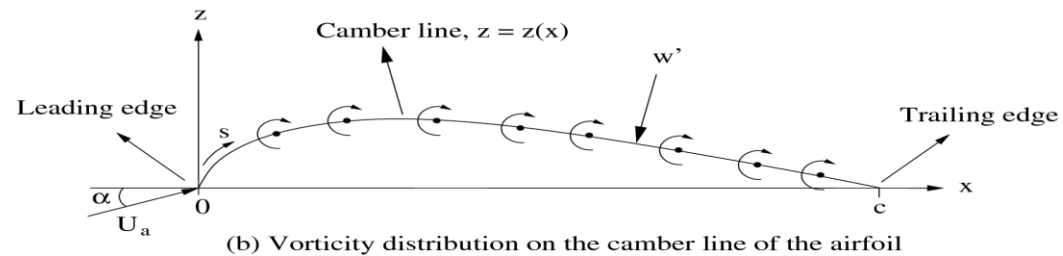
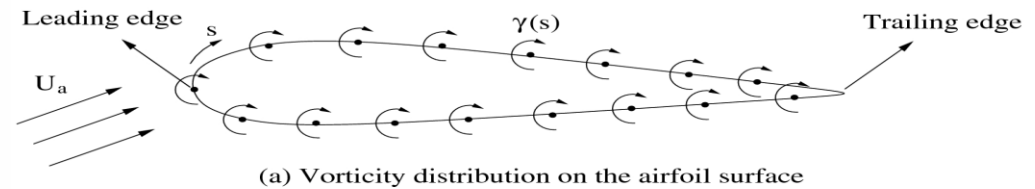
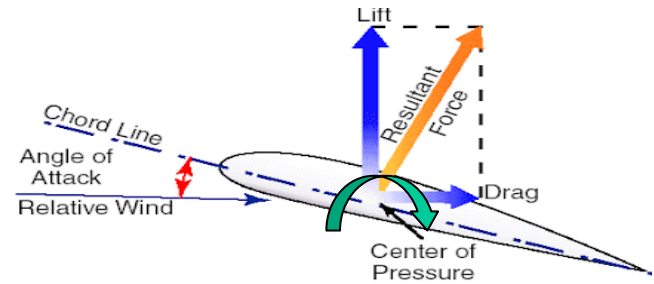
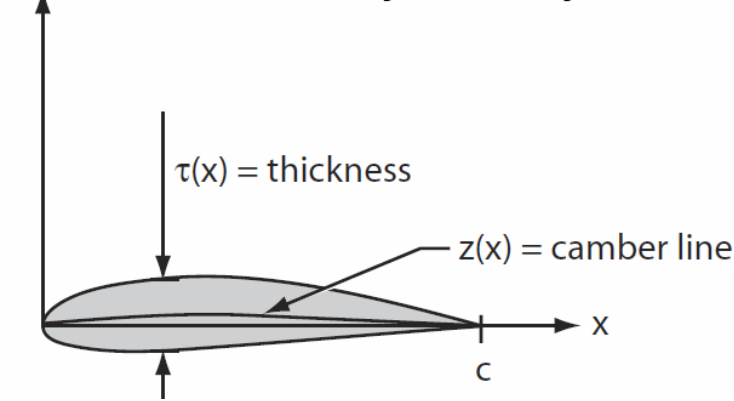
### Thin Airfoil Theory Summary



# Thin Airfoil Theory

- Replace thin airfoil with a camber line (assume small thickness and camber)
- Derivation of formula to predict the **aerodynamic force** and **moment** generated by a thin airfoil.
- Determination of the **pressure and aerodynamic centers**.
- Application of the thin airfoil theory to symmetric airfoils and cambered airfoils.

## Thin Airfoil Theory Summary

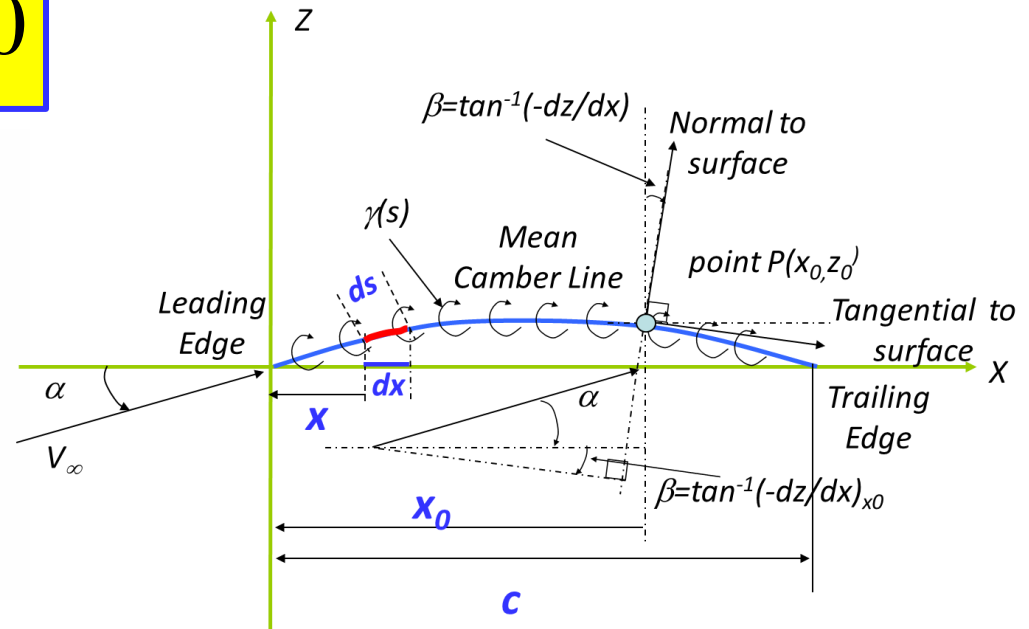
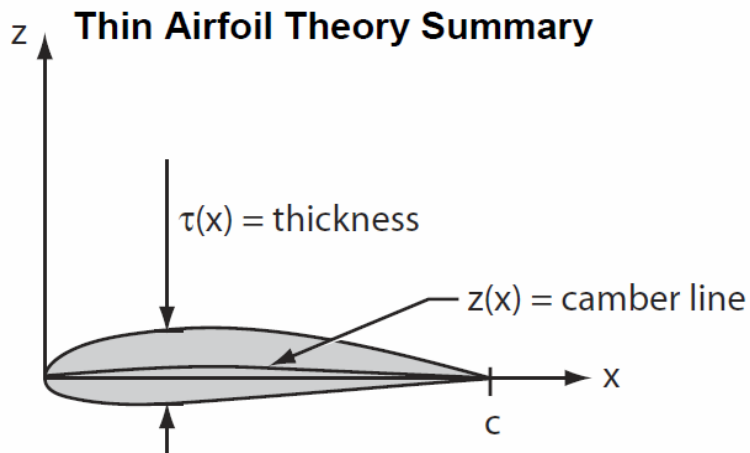


# Thin Airfoil Theory

## Principle:

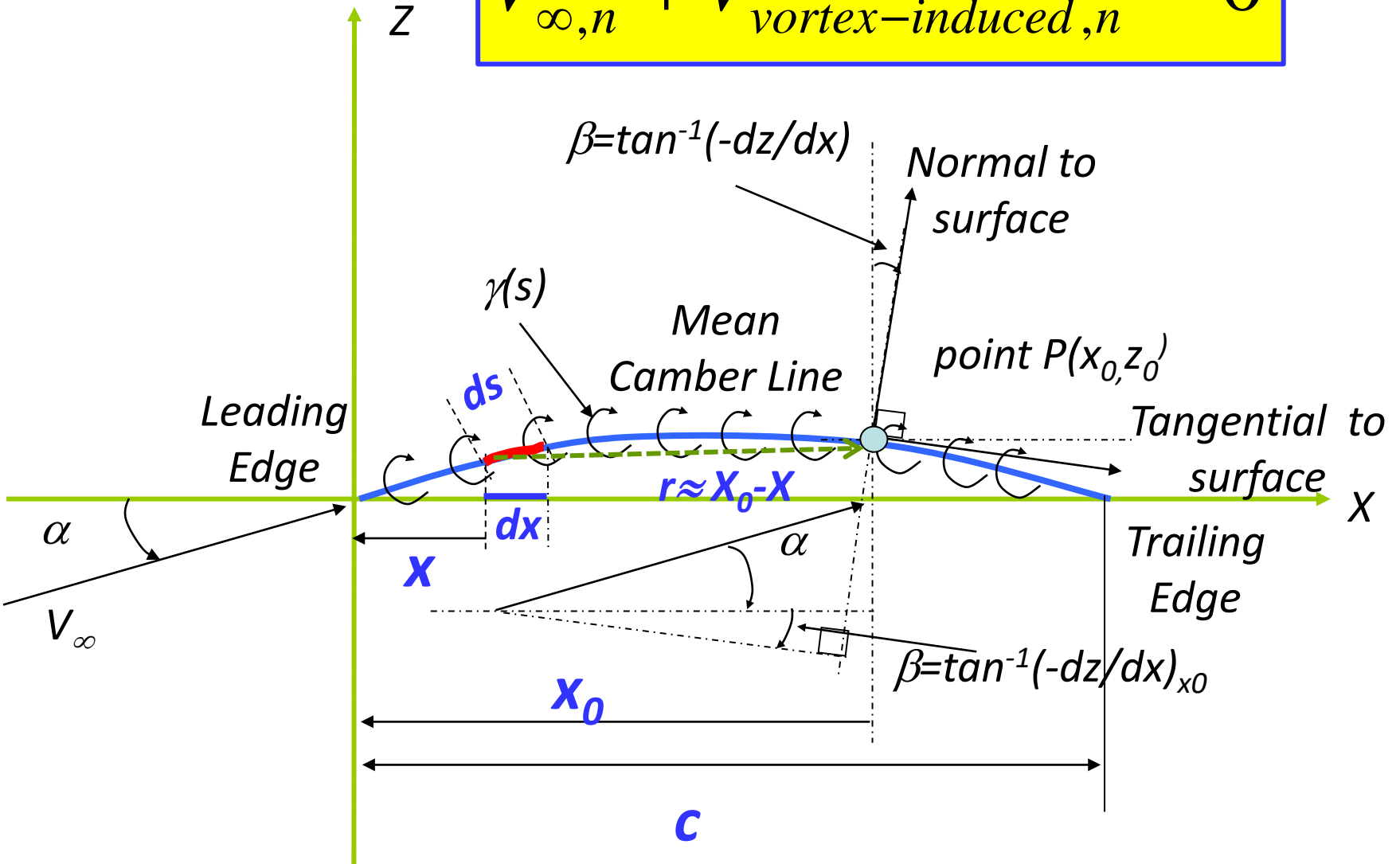
- Replace thin airfoil with the mean camber line (MCL) because of the small thickness and camber of the airfoil
- MCL assumed to be a streamline of the flow around the thin airfoil.
- To force the MCL to be a streamline, the sum of all velocity components normal to the MCL must be equal to zero.

$$V_{\infty, n} + V_{\text{vortex-induced}, n} = 0$$



# Thin Airfoil Theory

$$V_{\infty, n} + V_{\text{vortex-induced}, n} = 0$$



# Thin Airfoil Theory

- To force the mean camber line to be a streamline, the sum of all velocity components normal to the mcl must be equal to zero.
- Now determine the component of the freestream velocity normal to the mcl.

$$V_{\infty,n} = V_{\infty} \sin(\alpha + \varepsilon);$$

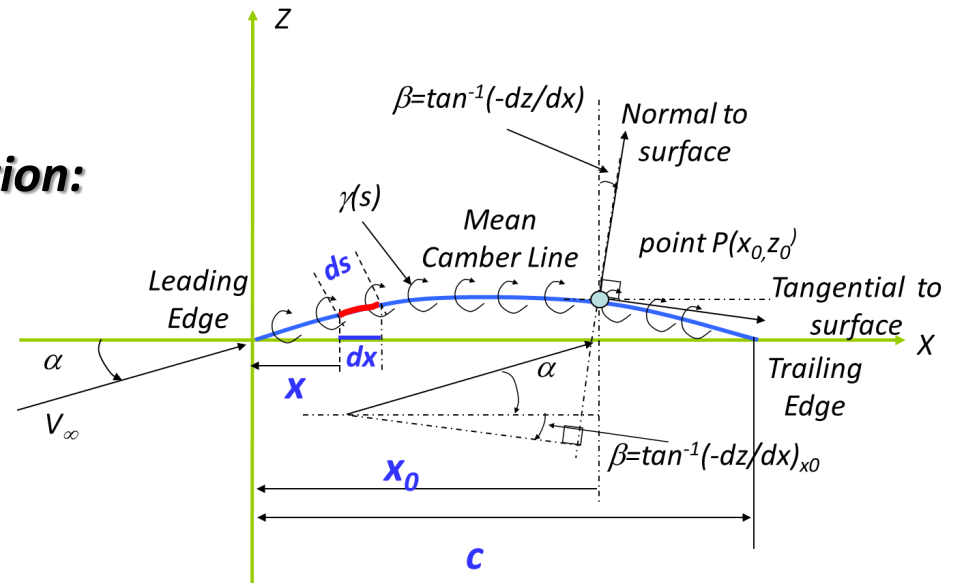
- Within thin airfoil theory approximation:**

Since the angle  $\varepsilon$  is small

$$\Rightarrow \beta = \tan^{-1}\left(-\frac{dz}{dx}\right) \approx -\frac{dz}{dx};$$

$$\text{Then: } V_{\infty,n} = V_{\infty} \sin(\alpha + \beta) = V_{\infty} \sin\left(\alpha - \frac{dz}{dx}\right);$$

Since the angle  $\alpha$  is small  $\Rightarrow \sin(\alpha + \beta) \approx \alpha + \beta \approx \alpha - \frac{dz}{dx};$



- Therefore:**

$$V_{\infty,n} = V_{\infty} \left(\alpha - \frac{dz}{dx}\right);$$

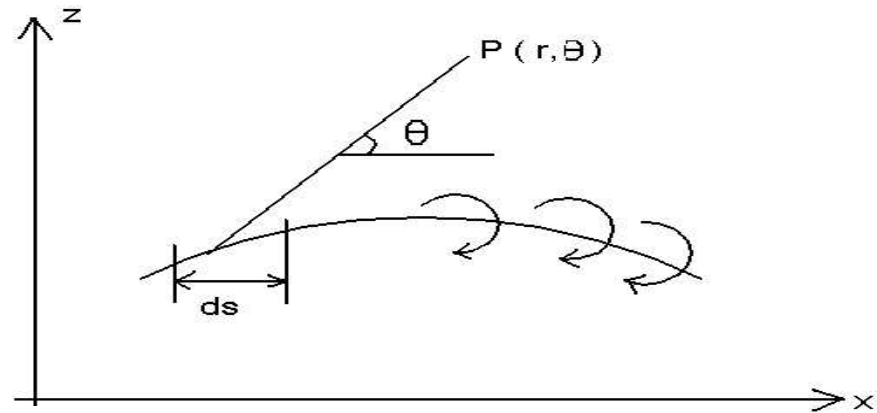
# Thin Airfoil Theory

- Consider the flow induced by an elemental vortex sheet  $ds$  at a point  $P$  on the vortex sheet.

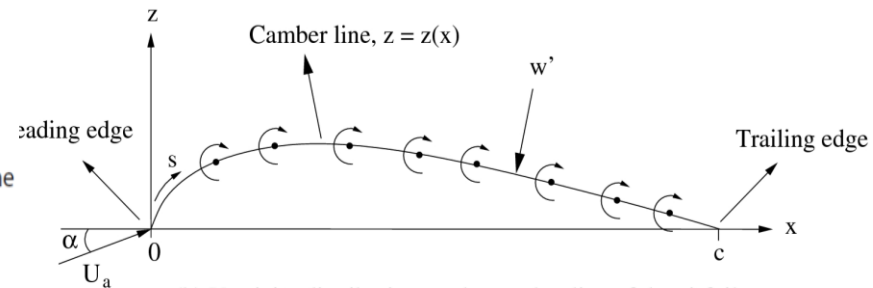
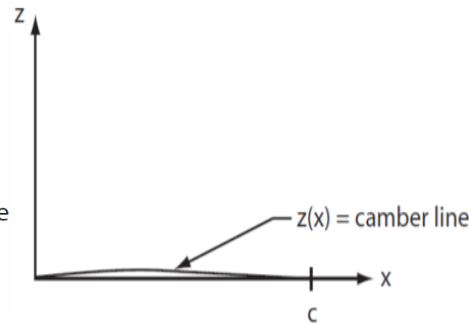
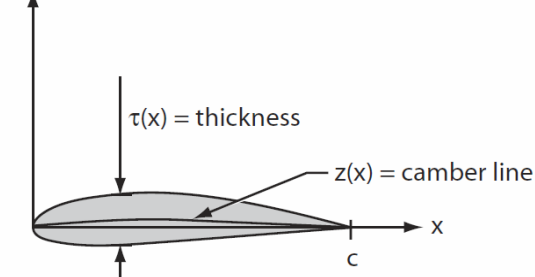
- Velocity induced by a 2-D vortex is: 
$$\vec{V} = V_\theta \hat{e}_\theta = -\frac{\Gamma}{2\pi r} \hat{e}_\theta; \quad .$$

- Similarly, the velocity at the point  $P$  induced by the vortex sheet of infinitesimal length  $ds$  is given by:

$$d\vec{V}_P = -\frac{\gamma(s)ds}{2\pi r} \hat{e}_\theta; \quad .$$



## Thin Airfoil Theory Summary



(b) Vorticity distribution on the camber line of the airfoil



# Thin Airfoil Theory

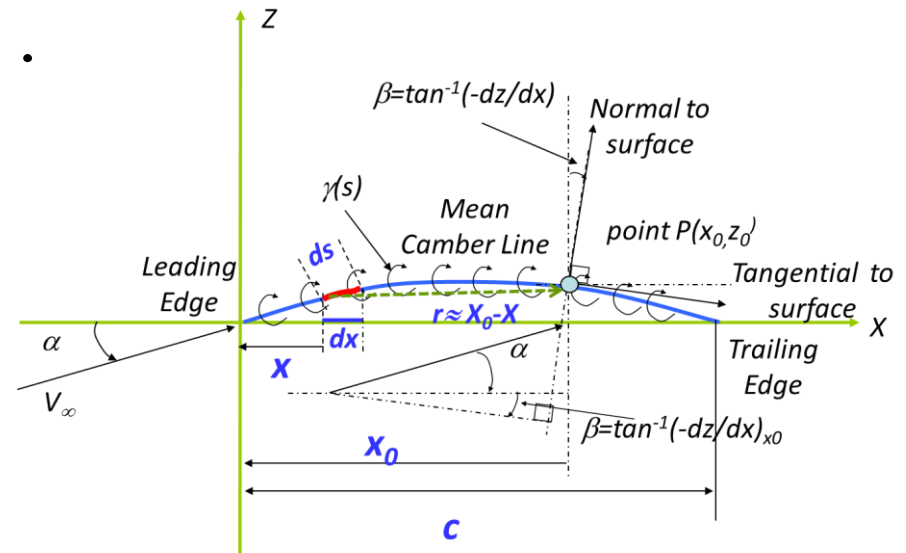
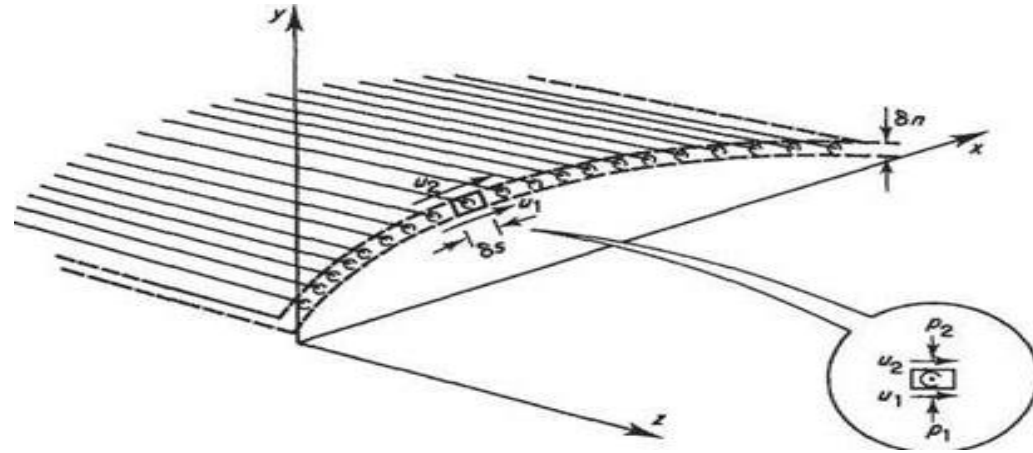
- Within thin airfoil theory approximation:

Thin airfoil  $\Rightarrow ds \approx dx$   
therefore :

$$d\vec{V}_P = -\frac{\gamma(s)ds}{2\pi r} \hat{e}_\theta;$$

$$\Rightarrow dV_{induced,n} = -\frac{\gamma(x)dx}{2\pi(x_0 - x)}$$

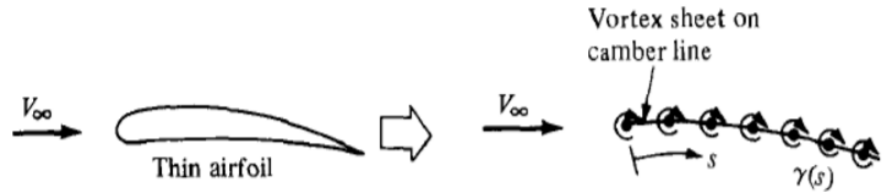
$$V_{vortex-induced,n} = -\int_0^c \frac{\gamma(x)dx}{2\pi(x_0 - x)}$$



# Thin Airfoil Theory

- To force the mean camber line to be a streamline, the sum of all velocity components normal to the mcl must be equal to zero.

$$V_{\infty,n} + V_{\text{vortex-induced},n} = 0$$



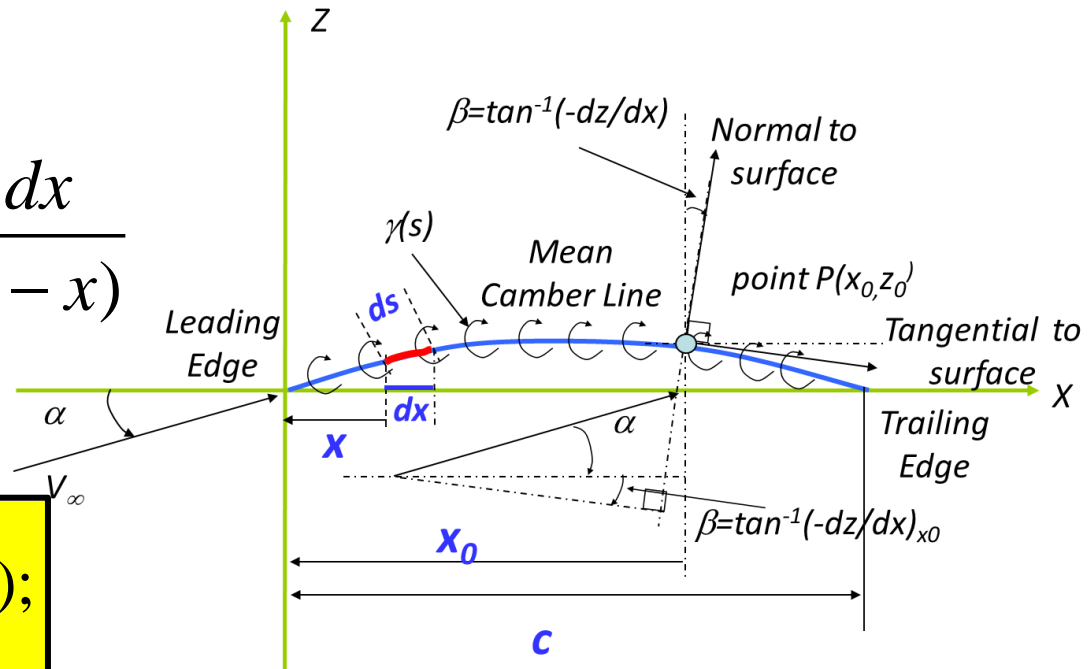
- Since

$$V_{\infty,n} = V_{\infty} \left( \alpha - \frac{dz}{dx} \right);$$

$$V_{\text{vortex-induced},n} = - \int_0^c \frac{\gamma(x) dx}{2\pi(x_0 - x)}$$

- Therefore:

$$\frac{1}{2\pi} \int_0^c \frac{\gamma(x) dx}{(x_0 - x)} = V_{\infty} \left( \alpha - \frac{dz}{dx} \right);$$



# Thin Airfoil Theory

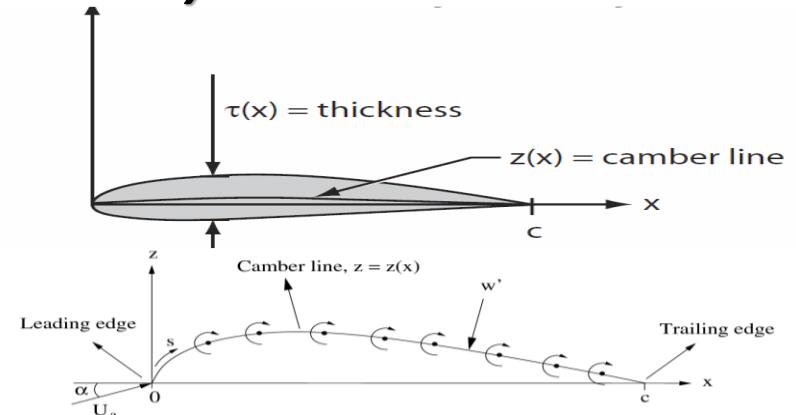
The integral equation of thin airfoil theory:

$$\frac{1}{2\pi} \int_0^c \frac{\gamma(x) dx}{(x_0 - x)} = V_\infty \left( \alpha - \frac{dz}{dx} \right);$$

For a given airfoil geometry,  $dz/dx$  is known. the only unknown in the above equation is  $\gamma(x)$ .

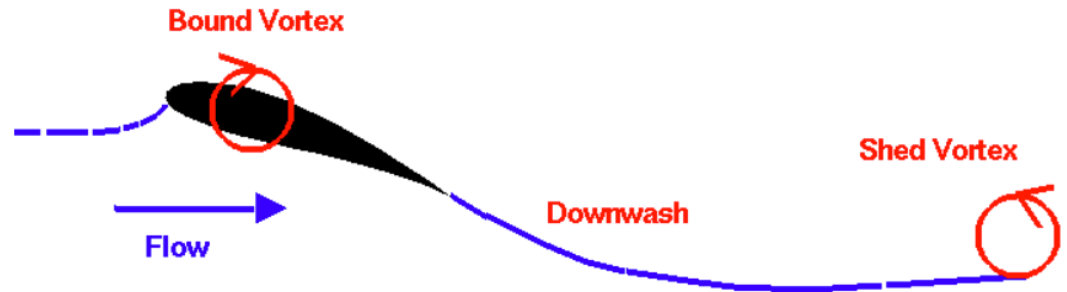
If  $\gamma(x)$  can be determined, then

$$\Gamma = \int_0^c \gamma(x) dx;$$



The Kutta-Joukowski Lift Theorem states the lift per unit length of an airfoil is equal to the density ( $\rho$ ) of the air times the strength of the rotation ( $\Gamma$ ) times the velocity ( $V$ ) of the air.

$$L' = \rho V_\infty \Gamma$$



# Thin Airfoil Theory

## The integral equation of thin airfoil theory:

$$\frac{1}{2\pi} \int_0^c \frac{\gamma(x) dx}{(x_0 - x)} = V_\infty \left( \alpha - \frac{dz}{dx} \right);$$

To solve the integral equation, we first make a transformation:

$$x = \frac{c}{2}(1 - \cos \theta); \quad \theta = (0, \pi)$$

If  $\theta = 0 \Rightarrow x = 0$ , at Leading Edge

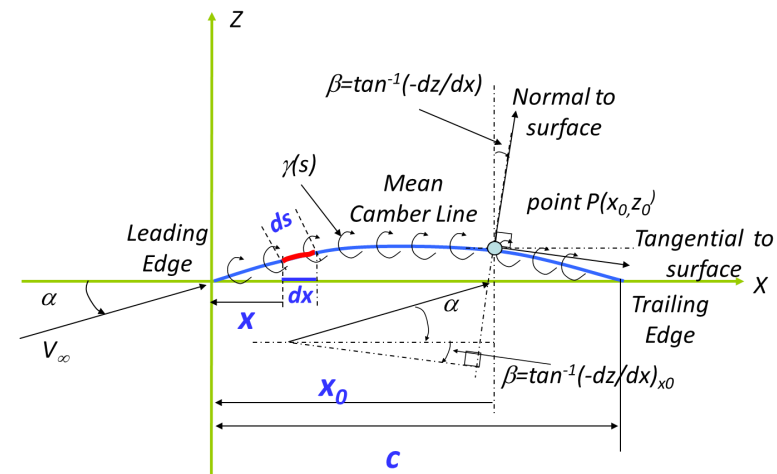
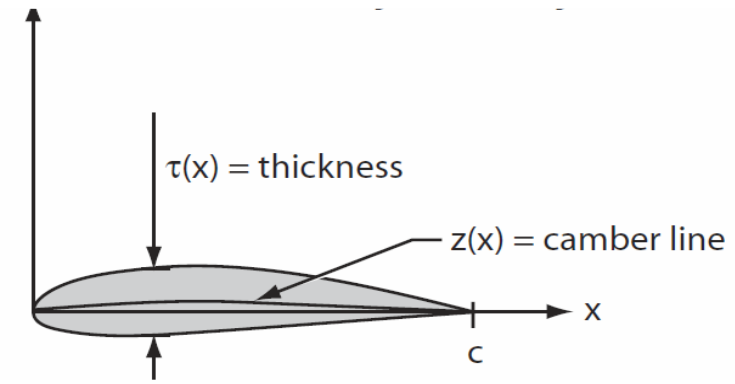
If  $\theta = \pi \Rightarrow x = c$ , at Trailing Edge

Then,  $x$  is a point on  $x_0$ , corresponding to  $\theta_0$ :

$$\text{Then: } x = \frac{c}{2}(1 - \cos \theta_0);$$

$$\frac{1}{2\pi} \int_0^\pi \frac{\gamma(\theta) \cdot c \cdot \sin \theta d\theta}{\frac{c}{2}(1 - \cos \theta_0) - \frac{c}{2}(1 - \cos \theta)} = V_\infty \left( \alpha - \frac{dz}{dx} \right);$$

$$\Rightarrow \frac{1}{2\pi} \int_0^\pi \frac{\gamma(\theta) \sin \theta d\theta}{\cos \theta - \cos \theta_0} = V_\infty \left( \alpha - \frac{dz}{dx} \right);$$



# Thin Airfoil Theory

The integral equation of thin airfoil theory:

$$\frac{1}{2\pi} \int_0^\pi \frac{\gamma(\theta) \sin \theta d\theta}{\cos \theta - \cos \theta_0} = V_\infty \left( \alpha - \frac{dz}{dx} \right);$$

Solution of the integral equation for a symmetrical airfoil:

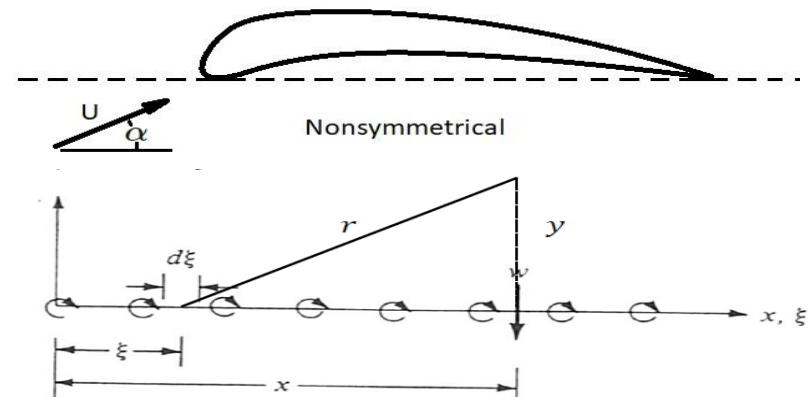
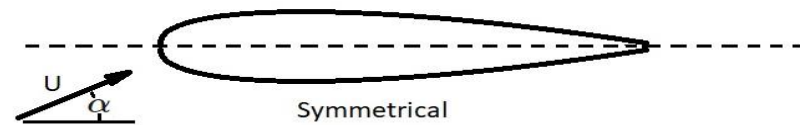
Since it is symmetrical airfoil, therefore:

$$\frac{dz}{dx} = 0$$

$$\frac{dz}{dx} = 0 \Rightarrow \frac{1}{2\pi} \int_0^\pi \frac{\gamma(\theta) \sin \theta d\theta}{\cos \theta - \cos \theta_0} = V_\infty \alpha;$$

The solution will be:

$$\gamma(\theta) = 2\alpha V_\infty \frac{1 + \cos \theta}{\sin \theta}$$



# Thin Airfoil Theory

□ **The integral equation of thin symmetrical airfoil.**

$$\frac{1}{2\pi} \int_0^\pi \frac{\gamma(\theta) \sin \theta d\theta}{\cos \theta - \cos \theta_0} = V_\infty \alpha;$$

□ **Solution of the equations:**

$$\gamma(\theta) = 2\alpha V_\infty \frac{1 + \cos \theta}{\sin \theta}$$

□ **Verification of the solution:**

$$\frac{1}{2\pi} \int_0^\pi \frac{\gamma(\theta) \sin \theta d\theta}{\cos \theta - \cos \theta_0} = \frac{1}{2\pi} \int_0^\pi \frac{2\alpha V_\infty \frac{1 + \cos \theta}{\sin \theta} \sin \theta d\theta}{\cos \theta - \cos \theta_0} = \frac{\alpha V_\infty}{\pi} \int_0^\pi \frac{(1 + \cos \theta) d\theta}{\cos \theta - \cos \theta_0}$$

$$\therefore \int_0^\pi \frac{\cos(n\theta) d\theta}{\cos \theta - \cos \theta_0} = \frac{\pi \sin(n\theta_0)}{\sin \theta_0}$$

$$\therefore \frac{\alpha V_\infty}{\pi} \int_0^\pi \frac{(1 + \cos \theta) d\theta}{\cos \theta - \cos \theta_0} = \frac{\alpha V_\infty}{\pi} \left[ \int_0^\pi \frac{d\theta}{\cos \theta - \cos \theta_0} + \int_0^\pi \frac{\cos \theta d\theta}{\cos \theta - \cos \theta_0} \right]$$

$$= \frac{\alpha V_\infty}{\pi} \left[ \frac{\pi \sin(0 \cdot \theta_0)}{\sin \theta_0} + \frac{\pi \sin(1 \cdot \theta_0)}{\sin \theta_0} \right] = \frac{\alpha V_\infty}{\pi} [0 + \pi] = \alpha V_\infty$$

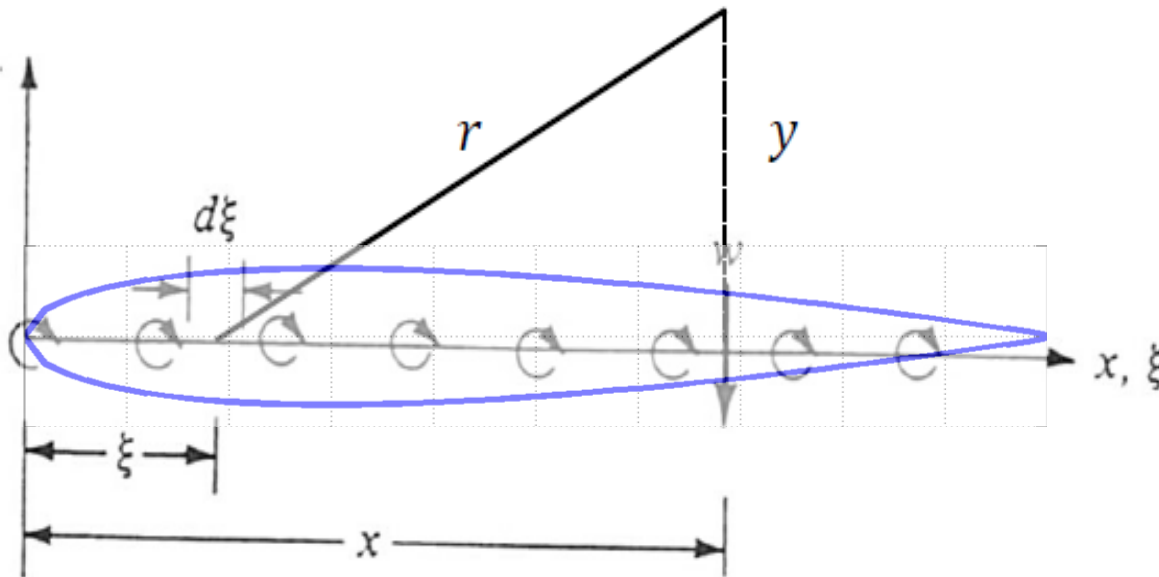
# Thin Airfoil Theory

- The integral equation of thin airfoil theory for symmetric airfoils:

$$\frac{1}{2\pi} \int_0^\pi \frac{\gamma(\theta) \sin \theta d\theta}{\cos \theta - \cos \theta_0} = V_\infty \alpha;$$

- Solution of the equation:

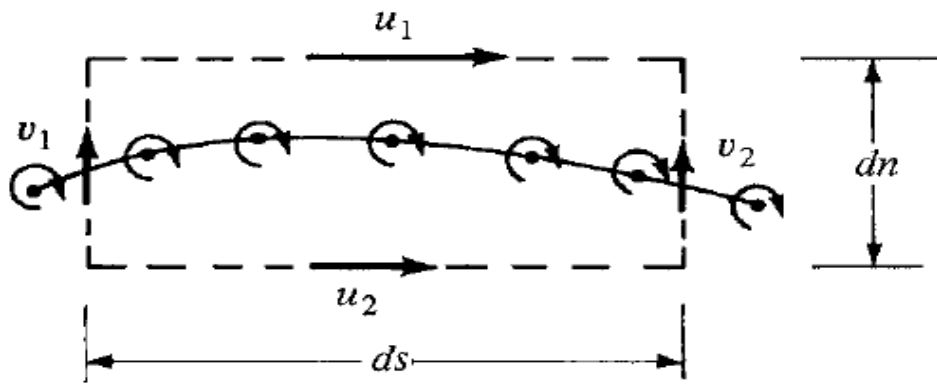
$$\gamma(\theta) = 2\alpha V_\infty \frac{1 + \cos \theta}{\sin \theta}$$



# Thin Airfoil Theory

## Kutta condition:

- For a given airfoil at a given angle of attack, the value of  $\Gamma$  around the airfoil is such that the flow would leave the trailing edge smoothly.

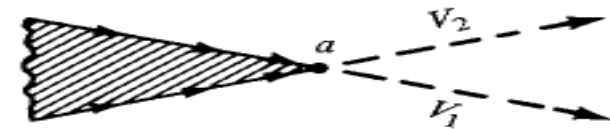
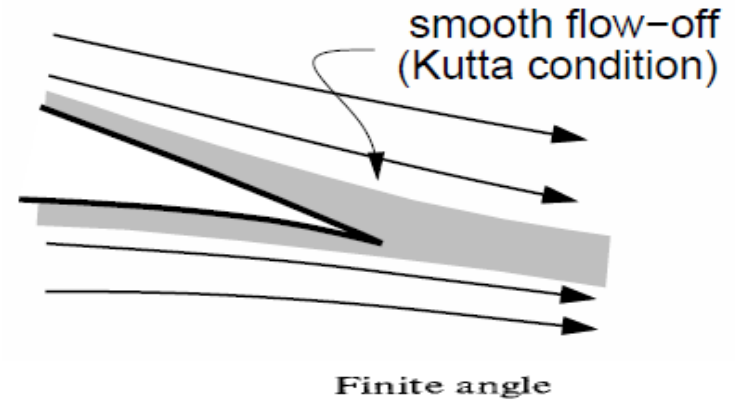


$$\gamma = u_2 - u_1$$

the vortex sheet

$$\gamma(\text{TE}) = V_2 - V_1$$

$$\gamma(\text{TE}) = 0$$



At point a:  $V_1 = V_2 = 0$

• Case #1

Cusp



At point a:  $V_1 = V_2 \neq 0$

• Case #2



# Thin Airfoil Theory

## Is Kutta condition satisfied at TE?

$$\gamma(\theta) = 2\alpha V_\infty \frac{1 + \cos \theta}{\sin \theta}$$

at airfoil LE:  $\theta = \pi$

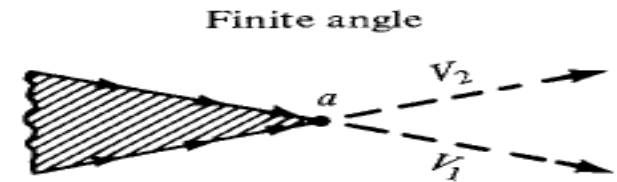
$$\Rightarrow \gamma(\pi) = 2\alpha V_\infty \frac{1 + \cos \pi}{\sin \pi} = \frac{0}{0}$$

Therefore:

$$\gamma(\theta) \Big|_{\theta \rightarrow \pi} = \frac{\frac{d(1 + \cos \theta)}{d\theta}}{\frac{d(\sin \theta)}{d\theta}} \Big|_{\theta \rightarrow \pi}$$

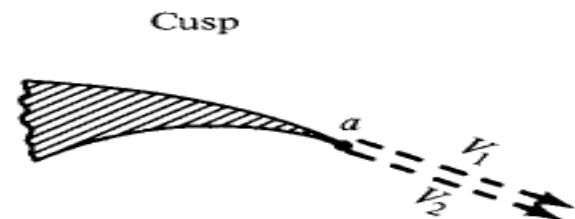
$$= \frac{-\sin \theta}{\cos \theta} \Big|_{\theta \rightarrow \pi} = \frac{0}{-1} = 0$$

$\Rightarrow$  Kutta condition is satisfied at TE!



At point  $a$ :  $V_1 = V_2 = 0$

### • Case #1



At point  $a$ :  $V_1 = V_2 \neq 0$

### • Case #2

- In relation to the vortex sheet discontinuity

$$\gamma(\text{TE}) = V_2 - V_1$$

$$\gamma(\text{TE}) = 0$$