

# AerE310: Incompressible Aerodynamics

## Homework Problem Set #1:

Due: 5:00 PM, Friday, 02/02/2024

1. Expand following terms:

a).  $\frac{d}{dt}(\vec{A} \cdot (\vec{B} \times \vec{C}))$

b).  $\frac{d}{dt}(\vec{A} \times (\vec{B} \times \vec{C}))$

**Solution:**

**a.**

$$\begin{aligned}\frac{d}{dt}(\vec{A} \cdot (\vec{B} \times \vec{C})) &= \frac{d\vec{A}}{dt} \cdot (\vec{B} \times \vec{C}) + \vec{A} \cdot \frac{d}{dt}(\vec{B} \times \vec{C}) \\ &= \frac{d\vec{A}}{dt} \cdot (\vec{B} \times \vec{C}) + \vec{A} \cdot \left( \frac{d\vec{B}}{dt} \times \vec{C} \right) + \vec{A} \cdot \left( \vec{B} \times \frac{d\vec{C}}{dt} \right)\end{aligned}$$

**b.**

$$\frac{d}{dt}(\vec{A} \times (\vec{B} \times \vec{C})) = \frac{d\vec{A}}{dt} \times (\vec{B} \times \vec{C}) + \vec{A} \times \left( \frac{d\vec{B}}{dt} \times \vec{C} \right) + \vec{A} \times \left( \vec{B} \times \frac{d\vec{C}}{dt} \right)$$

3. Find  $\nabla\Phi$  if

a).  $\Phi = \ln|\vec{r}|$

b).  $\Phi = \frac{1}{r}$

(Hint:  $\vec{r} = x\hat{i} + y\hat{j} + z\hat{k}$  in Cartesian coordinate system)

**Solution:**

a).

$$\Phi = \ln|\vec{r}| = \frac{1}{2} \ln(x^2 + y^2 + z^2)$$

$$\begin{aligned}\nabla\Phi &= \nabla\left[\frac{1}{2} \ln(x^2 + y^2 + z^2)\right] = \frac{x\hat{i} + y\hat{j} + z\hat{k}}{x^2 + y^2 + z^2} \\ &= \frac{\vec{r}}{r^2}\end{aligned}$$

b).

$$\Phi = \frac{1}{r} = \frac{1}{\sqrt{x^2 + y^2 + z^2}}$$

$$\begin{aligned}\nabla\Phi &= \nabla\left(\frac{1}{\sqrt{x^2 + y^2 + z^2}}\right) = -\frac{x\hat{i} + y\hat{j} + z\hat{k}}{(x^2 + y^2 + z^2)^{\frac{3}{2}}} \\ &= -\frac{\vec{r}}{r^3}\end{aligned}$$

4. Find directional derivative of  $\Phi = x^2yz + 4xz^2$  at point  $(1, -2, 1)$  in the direction of  $2\hat{i} - \hat{j} - 2\hat{k}$ .

**Solution:**

$$\phi = x^2yz + 4xz^2 \quad \nabla\phi = (2xyz + 4z^2, x^2z, x^2y + 8xz)$$

at point  $(1, -2, 1)$ ,  $\nabla\phi = (0, 1, 6)$  unit vector for  $(2, -1, -2)$  is  $(\frac{2}{3}, -\frac{1}{3}, -\frac{2}{3})$  the directional derivative is  $(0, 1, 6) \cdot (\frac{2}{3}, -\frac{1}{3}, -\frac{2}{3}) = -4\frac{1}{3}$

5. If  $\vec{R} = \vec{R}(t) = r\hat{e}_r + z\hat{e}_z$  is the position vector of a particle in cylindrical coordinates, Obtain expression for velocity vector,  $\vec{V}$ , and acceleration vector,  $\vec{a}$ , at that point.

**Solution:**

$$\because \vec{R} = r\hat{e}_r + z\hat{e}_z$$

$$\therefore \vec{V} = \frac{d\vec{R}}{dt} = \frac{d(r\hat{e}_r + z\hat{e}_z)}{dt} = \frac{dr}{dt}\hat{e}_r + r\frac{d\hat{e}_r}{dt} + \frac{dz}{dt}\hat{e}_z + z\frac{d\hat{e}_z}{dt}$$

$\because \hat{e}_r, \hat{e}_z, \hat{e}_\theta$  is the functions of the position  $(r, \theta, z)$ , and the position  $(r, \theta, z)$  are the functions of time,  $t$ .

$\therefore$

$$\frac{d\hat{e}_r}{dt} = \frac{\partial \hat{e}_r}{\partial r} \frac{\partial r}{\partial t} + \frac{\partial \hat{e}_r}{\partial \theta} \frac{\partial \theta}{\partial t} + \frac{\partial \hat{e}_r}{\partial z} \frac{\partial z}{\partial t} = 0 + \hat{e}_\theta \frac{\partial \theta}{\partial t} + 0 = \hat{e}_\theta \frac{\partial \theta}{\partial t}$$

$$\frac{d\hat{e}_\theta}{dt} = \frac{\partial \hat{e}_\theta}{\partial r} \frac{\partial r}{\partial t} + \frac{\partial \hat{e}_\theta}{\partial \theta} \frac{\partial \theta}{\partial t} + \frac{\partial \hat{e}_\theta}{\partial z} \frac{\partial z}{\partial t} = 0 + (-\hat{e}_r) \frac{\partial \theta}{\partial t} + 0 = -\hat{e}_r \frac{\partial \theta}{\partial t}$$

$$\frac{d\hat{e}_z}{dt} = \frac{\partial \hat{e}_z}{\partial r} \frac{\partial r}{\partial t} + \frac{\partial \hat{e}_z}{\partial \theta} \frac{\partial \theta}{\partial t} + \frac{\partial \hat{e}_z}{\partial z} \frac{\partial z}{\partial t} = 0 + 0 + 0 = 0$$

Defining:  $V_r = \frac{dr}{dt}$ ;  $V_\theta = \frac{d\theta}{dt}$ ;  $V_z = \frac{dz}{dt}$  Then:

$$\begin{aligned} \vec{V} &= \frac{d\vec{R}}{dt} = \frac{dr}{dt}\hat{e}_r + r\frac{d\hat{e}_r}{dt} + \frac{dz}{dt}\hat{e}_z + z\frac{d\hat{e}_z}{dt} \\ &= V_r\hat{e}_r + rV_\theta\hat{e}_\theta + V_z\hat{e}_z \end{aligned}$$

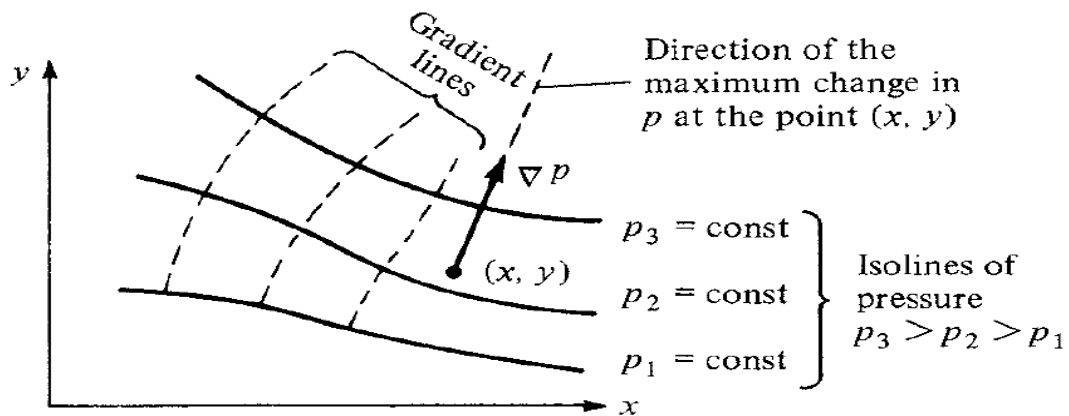
Defining:  $a_r = \frac{dV_r}{dt}$ ;  $a_\theta = \frac{dV_\theta}{dt}$ ;  $a_z = \frac{dV_z}{dt}$  Acceleration vector:

$$\begin{aligned} \vec{a} &= \frac{d\vec{V}}{dt} = \frac{dV_r}{dt}\hat{e}_r + V_r\frac{d\hat{e}_r}{dt} + \frac{dV_\theta}{dt}\hat{e}_\theta + V_\theta\frac{d\hat{e}_\theta}{dt} + \frac{dV_z}{dt}\hat{e}_z + V_z\frac{d\hat{e}_z}{dt} \\ &= a_r\hat{e}_r + rV_rV_\theta\hat{e}_\theta + V_rV_\theta\hat{e}_\theta + ra_\theta\hat{e}_\theta + rV_\theta(-\hat{e}_rV_\theta) + a_z\hat{e}_z + 0 \\ &= [a_r - rV_\theta^2]\hat{e}_r + [2rV_rV_\theta + ra_\theta]\hat{e}_\theta + a_z\hat{e}_z \end{aligned}$$

6. Show that the directions of the isoline and the gradient line at any given points in a scalar field are orthogonal to each other. (Hint: use the concept of directional derivative)

**Solution:**

Consider an isoline with the scalar  $\phi = \text{const}$ . If the direction of the isoline be denoted by "s" at a point, then the directional derivative of  $\phi$  in the direction of "s" is given by:  $\frac{d\phi}{ds} = \nabla\phi \cdot \hat{e}_s$  where  $\hat{e}_s$  is a unit vector in the direction of  $\hat{s}$   $\frac{d\phi}{ds} = 0$ , since  $\phi$  does not vary on the isoline along which  $ds$  is taken.  $\therefore \nabla\phi \cdot \hat{e} = 0$ . Therefore  $\nabla\phi$  and  $\hat{e}$  are orthogonal.



7. Spherical coordinate  $(R, \varphi, \theta)$  are defined by the following inverse transformation:

$$\begin{aligned}x &= (R \sin \varphi) \cos \theta \\y &= (R \sin \varphi) \sin \theta \\z &= R \cos \varphi\end{aligned}$$

Where

$$\begin{aligned}0 &\leq R \leq \infty \\0 &\leq \theta \leq 2\pi \\0 &\leq \varphi \leq 2\pi\end{aligned}$$

- Obtain the scale factors for the spherical coordinate system.
- Obtain the unit vectors in spherical system as the function of Cartesian unit vectors.
- Obtain the derivatives of the unit vectors with respect to spherical coordinate directions and simplify the results to be only functions of spherical coordinates.
- Using vector algebra to obtain the divergence of a general vector in spherical coordinates. Simplify the results to be in conservation form.

**Solution**

**a.**

$$h_k = 1 \quad h_\varphi = R \quad h_\theta = k \sin \varphi$$

**b.**

$$\begin{aligned}\hat{e}_R &= \sin \varphi \cos \theta \hat{i} + \sin \varphi \sin \theta \hat{j} + \cos \varphi \hat{k} \\ \hat{e}_\varphi &= \cos \varphi \cos \theta \hat{i} + \cos \varphi \sin \theta \hat{j} - \sin \varphi \hat{k} \\ \hat{e}_\theta &= -\sin \theta \hat{i} + \cos \theta \hat{j}\end{aligned}$$

**c.**

$$\begin{aligned}\frac{\partial \hat{e}_n}{\partial k} &\equiv 0 & \frac{\partial \hat{e}_n}{\partial \varphi} &\equiv \hat{e}_\varphi & \frac{\partial \hat{e}_R}{\partial \theta} &\equiv \sin \varphi \hat{e}_\theta \\ \frac{\partial \hat{e}_\varphi}{\partial R} &\equiv 0 & \frac{\partial \hat{e}_\varphi}{\partial \varphi} &\equiv -\hat{e}_k & \frac{\partial \hat{e}_\varphi}{\partial \theta} &\equiv \cos \varphi \hat{e}_\theta \\ \frac{\partial \hat{e}_\theta}{\partial R} &\equiv 0 & \frac{\partial \hat{e}_\theta}{\partial \varphi} &\equiv 0 & \frac{\partial \hat{e}_\theta}{\partial \theta} &= -(\sin \varphi \hat{e}_k + \cos \varphi \hat{e}_\varphi)\end{aligned}$$

**d.**

$$\begin{aligned}\nabla \cdot \vec{v} &= \left( \hat{e}_n \frac{\partial}{\partial R} + \frac{\hat{e}_\varphi}{R} \frac{\partial}{\partial \varphi} + \frac{\hat{e}_\theta}{R \sin \varphi} \frac{\partial}{\partial \theta} \right) \cdot (V_R \hat{e}_k + V_\varphi \hat{e}_\varphi + V_\theta \hat{e}_\theta) \\ &= \hat{e}_R \frac{\partial}{\partial R} \cdot (V_R \hat{e}_k + V_\varphi \hat{e}_\varphi + V_\theta \hat{e}_\theta) + \frac{\hat{e}_\varphi}{R} \frac{\partial}{\partial \varphi} \cdot (v_R \hat{e}_k + v_\varphi \hat{e}_\varphi + v_\theta \hat{e}_\theta) + \frac{\hat{e}_\theta}{k \sin \varphi} \frac{\partial}{\partial \varphi} \cdot [W_R \hat{e}_n + V_\varphi \hat{e}_\varphi + v_\theta \hat{e}_\theta] \\ &= \left( \frac{\partial V_k}{\partial R} + \frac{V_k}{R} + \frac{V_R}{R} \right) + \left( \frac{1}{R} \frac{\partial V_\varphi}{\partial \varphi} + \frac{V_\varphi \cos \varphi}{R \sin \varphi} \right) + \left( \frac{1}{R \sin \varphi} \frac{\partial V_\theta}{\partial \theta} \right) \\ &= \frac{1}{R^2} \left( \frac{\partial R^2 V_R}{\partial R} \right) + \frac{1}{R \sin \varphi} \frac{\partial}{\partial \varphi} (V_\varphi \sin \varphi) + \frac{1}{R \sin \varphi} \frac{\partial V_\theta}{\partial \theta} \\ &= \frac{1}{R^2 \sin \varphi} \left[ \frac{\partial (R^2 \sin \varphi V_R)}{\partial R} + \frac{\partial (V_\varphi \sin \varphi)}{\partial \varphi} + \frac{\partial (R V_\theta)}{\partial \theta} \right]\end{aligned}$$

8. Find the acceleration of a fluid particle at  $(r, \theta, z)$  in cylindrical coordinate system.

**Solution:**

$$\begin{aligned}\vec{a} &= \frac{\partial \vec{v}}{\partial t} + \vec{v} \cdot \nabla \vec{v} \\ \vec{v} \cdot \nabla &= (v_r \hat{e}_r + v_\theta \hat{e}_\theta + v_z \hat{e}_z) \cdot \left( \hat{e}_r \frac{\partial}{\partial R} + \frac{\hat{e}_\theta}{R} \frac{\partial}{\partial \theta} + \hat{e}_z \frac{\partial}{\partial z} \right) \\ &= V_R \frac{\partial}{\partial R} + V_\theta \frac{\partial}{\partial \theta} + V_z \frac{\partial}{\partial z} \\ \vec{v} \cdot \nabla \vec{v} &= \left( V_R \frac{\partial}{\partial R} + \frac{V_\theta}{R} \frac{\partial}{\partial \theta} + V_z \frac{\partial}{\partial z} \right) (V_R \hat{e}_r + V_\theta \hat{e}_\theta + v_z \hat{e}_z) \\ &= V_R \frac{\partial V}{\partial R} \hat{e}_r + V_R \frac{\partial V_\theta}{\partial R} \hat{e}_\theta + V_R \frac{\partial V_z}{\partial R} \hat{e}_z + 0 + 0 + 0 \\ &\quad + \frac{V_\theta}{R} \frac{\partial V_R}{\partial \theta} \hat{e}_r + \frac{V_\theta V_R}{R} \hat{e}_\theta + \frac{V_\theta}{R} \frac{\partial V_\theta}{\partial \theta} \hat{e}_\theta + \frac{V_\theta^2}{R} (-\hat{e}_\theta) \\ &\quad + \frac{V_\theta}{R} \frac{\partial V_z}{\partial \theta} \hat{e}_z + 0 + V_z \frac{\partial V_r}{\partial z} \hat{e}_r + v_z \frac{\partial v_\theta}{\partial z} \hat{e}_\theta + V_z \frac{\partial V_z}{\partial z} \hat{e}_z + 0 + 0 + 0\end{aligned}$$

Using the derivatives of the unit vectors.

$$\begin{aligned}a_R &= \frac{\partial V_R}{\partial t} + V_R \frac{\partial V_R}{\partial R} + \frac{V_\theta}{R} \frac{\partial V_R}{\partial \theta} + V_z \frac{\partial V_R}{\partial z} - \frac{V_\theta^2}{R} \\ a_\theta &= \frac{\partial V_\theta}{\partial t} + V_R \frac{\partial V_\theta}{\partial R} + \frac{V_\theta V_R}{R} + \frac{V_\theta}{R} \frac{\partial V_\theta}{\partial \theta} + V_z \frac{\partial v_\theta}{\partial z} \\ a_z &= \frac{\partial v_z}{\partial t} + \psi k \frac{\partial V_z}{\partial k} + \frac{V_\theta \partial V_z}{R} + V_z \frac{\partial V_z}{\partial z}\end{aligned}$$